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## RESEARCH MEMORANDUM

AN INVESTIGATION OF THE LOW-SPEED STABILITY AND CONTROL CHARACTERISTICS OF SWEPT-FORWARD AND SWEPT-BACK WINGS IN THE AMES 40- BY 80-FOOT WIND TUNNEL

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NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS

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WASHINGTON June 10, 1947

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### NATIONAL ADVISORY COMMITTEE FOR AFRONAUTICS

### RESEALCH MEMORANDUM

AN INVESTIGATION OF THE LOW-SPEED STABILITY AND CONTROL CHARACTERISTICS OF SWEFT-FORWARD AND SWEPT-BACK WINGS
IN THE AMES 40- BY SO-FOOT WIND TURNEL
By Gerald N. McCormack and Viotor I. Stevens, Jr.

### SUMMARY

An investigation has been made at large scale of the characteristics of highly swept wings. Data were obtained at several angles of sideslip on wings having angles of sweep of ±45°, ±30°, and 0°. The airfoil sections of the wings varied from approximately NACA 0015 at the root to NACA 23009 at the tip. Each wing was investigated with flaps undeflected, partial-span split flaps deflected 60°, full-span split flaps deflected 60° and split-flap-type allerons deflected ±15°. Values of maximum lift were obtained at Reynolds numbers ranging from 5.7 to 9.2×10°. In this report the summarized results are compared with the predictions made by use of the simplified theory for the effect of sweep and with existing small-scale data. The basic wind-turnel results from which these summary data were taken are included in an appendix.

The primary problems accompanying the use of sweep as revealed by this investigation are the loss in maximum lift, the high effective dihedral, and the sharp reduction in lateral—control effectivenese. In general, eimple theory enables good predictions to be made of the gross effects of sweep but further

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refinements are necessary to obtain the accuracy required for design purposes. In cases where comparisons can be made, the indications are that, as sweep increases, scale effects diminish and large-scale results approach small-scale results.

### **INTRODUCTION**

Theory indicates and experiment has shown that the prime aerodynamic effect of wing sweep is to reduce by the essine of the angle of sweep the effective flight velocity experienced by the airfoil sections of the wing. This then enables increases in maximum flight speed to be attained before serious compressibility effects are encountered. Theory and experiment also show that wing sweep introduces a number of stability and control problems, the seriousness of which becomes accentuated at low flight speeds.

Small-scale tests have pointed out the general nature of these problems and indicated those which must be overcome if the high-speed benefits of sweep are to be realized. They have also suggested that boundary-layer flow and, hence, Reynolds number has a profound influence on measured characteristics and that the value of small-scale tests remain somewhat doubtful until the extent of this influence is understood.

Since nolar re-scale data were available for wings with large angles of sweep, an investigation of the offects of sweep was conducted in the Ames 40- by 80-foot wind tunnel and the results are reported herein. It is believed that

these data will go far towards establishing the datum required to estimate the effects of scale on highly swept wing plan forms. With this knowledge at hand it is evident that the value of future small-scale tests will be considerably increased.

This report discusses a summary of the basic results and compares them with simple swept-wing theories and, where possible, with existing small-scale data (references 1, 2, and 3). To make the basic data available for further analyses they are included as an appendix to this report.

### DESCRIPTION OF HODELS

The five models tested were composed of wing panels-from an available airplane which were given the desired plan form and sweep by individually fabricated tips and center sections. The resulting angles of sweep were 0°, 30°, and 45° sweepforward, and 30°, and 45° sweepback (measured with reference to the quarter-chord line of the airfoil sections). Aside from the angle of sweep, the prime plan-form variable was considered to be aspect ratio. The tips and center sections were constructed to give the smallest variation of this parameter possible without modification of the cirplane wing banels. No special attempt was made to control the variation of taper ratio, area or span. Photographs of the wings and plan-form drawings with pertinent dimensions are shown in figures 1 and 2. The geometric digraoteristics of the five wings tested are listed in table 1.

The airfoil sections of the swept wings were dictated by the sections of the airplane wing panels (an NACA 0015 at the inboard end and an NACA 23009 at the outboard end of the panel). The profiles of the center sections and tips were simply extencions of the wing-panel airfoil. To expedite construction, three tips only were fabricated; one for the swept-forward wings, one for the straight wing, and one for the swept-back wings. Thus for the swept-forward and swept-back wings compromise tips were used which were misalined 7½0 to the air stream. The twist in the chord plane of the wing panels was approximately 1/40 of washout. The dihedral of the chord-plane leading edge was kept at 00.

No attempt was made to improve the fairness of the wing panels beyond the original manufacturing condition. Thus, due to presence of various access plates, panel joints, etc., the wings were rough to a greater degree than that normally associated with latest construction requirements.

Partial-span and full-span split flaps were tested on all models. The flaps were 0.70 chord and were deflected 60°. The span of the partial-span flaps was 0.623 wing span for all models; the epan of the full-span flaps varied slightly from full-span (in no cass more than 0.064 wing span) as shown in table I.

Ailerons were simulated by attaching the outboard portion of one of the flaps to the right wing and deflecting it ±15° (up-deflection was obtained by attaching the flap to the upper surface of the wing). Thus the allerons as tested were

Except where noted, all chords and spans used in this report were measured parallel and perpendicular to the plane of symmetry. Flap defisction angles were measured in a plane perpendicular to the flap hinge line.

CL

0.20-chord split-flap-type allerons.

The wings were mounted on a faired sting which in turn was attached to the three-strut support system. Photographs of the wing installations are shown in figure 1.

### COEFFICIENTS AND SYMBOLS

The data are presented in the form of standard NACA coefficients and symbols as defined in figure 3 and the following tabulation. All forces and moments are presented about the stability axes with their origin located on the root chord, or root chord projected and at the same fore and aft location as the quarter N.A.C.

lift coefficient (lift/qS)

CD.	drag coefficient (drag/q3)
C <sub>m</sub>	pitching-moment coefficient (pitching moment)
Cı	rolling-moment coefficient (rolling moment)
c <sub>n</sub>	yawing moment coefficient (waving moment)
C <sub>Y</sub>	side-force coefficient (side force)
CLa	rate of change of lift coefficient with angle of
	attack, per degree
ACL	increment of lift coefficient due to deflecting
	flaps
CLmax	maximum lift coefficient
Cle	rate of change of rolling-moment coefficient with
( En Je de in de	sideslip, per degree
Clo	rate of change of rolling-moment coefficient with
	Ving-tip helix engle, per radian

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Cnp	rate of change of yawing-moment coefficient with
	sideelip, per degree
Claa	rate of change of rolling-moment coefficient with
	mileron angle, per degree
901 6/901	rate of change of Ct with lift coefficient where
ACIE /ACT	rate of change of Cip with lift coefficient where
dCng/dCL	rate of change of Cng with the lift coefficient squared
	squared squared
L/D	ratio of lift to drag
q	dynamic pressure, pounds per square foot
V	velocity along flight path, feet per second
α	angle of attack, degrees
β	angle of sideslip, degrees
Δ	angle of sweep of quarter chord line of airfoil
	sections, degrees (Sweepback is positive and
	sweepforward is negative.)
e la la la	effective dihedral, degrees
	control surface deflection, degrees
Non-French	aspect ratio based on span (b2)
	aspect ratio based on length of quarter chord
	Tine (S con de )
	Jones' edge-velocity correction
	taper ratio, ratio of tip shord to root chord $\left(\frac{a_t}{a_r}\right)$
109.00	ving area, equara feet

### NACA RM No. A6K15

- o. mean serodynamic chord of wing measured parallel
  to plane of symmetry, feet
- b wing span measured perpendicular to the plane of symmetry, feet
- ci wing-tip chard
- cr wing-root chord

### TESTS AND RESULTS

For each of the model configurations six-component force and moment data were obtained through an angle-of-attack range at each of several angles of sideslip. The data were obtained at dynamic pressures which rarge from 5 to 75 pounds per square foot ( $R = 2.8 \times 10^6$  to  $R = 15.0 \times 10^6$ ) most of the data were obtained at dynamic pressures of 10 to 20 pounds per square foot ( $R = 4.0 \times 10^6$  and  $R = 5.3 \times 10^6$ , respectively).

The basic data obtained from the wind-tunnel tests of the five swept wings are described in the appendix. Also included in the appendix is a description of the corrections and tares applied to the data.

### DISCUSSION

In this discussion an evaluation is made of the effect of wing sweep on the more important aerodynamic parameters and of the consequent effect on sirplans performance and etability.

These Reynolds numbers are based upon the K.A.C. as a reference length and are the minimum and maximum limits of the variation including the change in chord length with sweep.

Also, the accouracy with which the simplified ewern theory may be used to predict the characteristics of swept wings is evaluated by a comparison with the experimental data. Finally, an attempt is made to compare at least qualitatively the values of the various characteristics as obtained at small-scale  $(R < 1.5 \times 10^6)$  and full-scale Reynolds numbers. The summary data on which this discussion is based have been extracted (for a test dynamic pressure of 20 lb/eq ft) from the measured characteristics included in the appendix.

The concepts advanced by Betz in reference 4 form the groundwork for the theory of the aerodynamic effects of incorporating eweep in a wing plan form. These concepts are based on the assumption that for an infinite-span wing only the velocity component normal to the quarter-chord line influences the pressures over a wing; the spanwise component of velocity is neglected. Thus, if the velocity components are resolved perpendicular and parallel to the quarter-chord line of a wing, the effective dynamic pressure over the wing will decrease in proportion to the square of the cosine of the angle of sweep and the effective angle of attack will increase in proportion to the reciprocal of the cosine of the angle of sweep. These changes in effective dynamic pressure and angle of attack brought about by wing sweep form the basis for the existing simplified eweep theory.

In interpreting the comparisons to be made between the simplified theory and experimental results, the limitations

of the simplified theory must be borne in mind. Over the root section of highly swept finite-span wings, particularly highly tapered low aspect ratio wings, the basic assumption that the wing reacts only to air velocities normal to the quarter-chord line probably does not hold. It should also be noted that simplified theory in its present form applies only to wings which generate an additional loading due to angle-of-attack change that is rectangular in form. Therefore appreciable deviations from rectangular loading such as produced by taper will result in discrepancies between the theoretical and experimental results.

### Lift Characteristics

Lift-curve slope. The simplified theory indicates a decrease in lift-curve slope proportional only to 60sA. To account for induction effects, a correction must also be made for any variations of aspect ratio. Hence, the effect of sweep on lift-curve slope, when corrected for aspect ratio, will be in accordance with the relation:

$$\left(c_{L_{\alpha}}\right)_{\Lambda} = \left(c_{L_{\alpha}}\right)_{\Lambda=0} \operatorname{cos}_{\Lambda} \frac{\left[\Lambda/(\Lambda+2)\right]_{\Lambda}}{\left[\Lambda/(\Lambda+2)\right]_{\Lambda}}$$

In conformity with standard nomenclature, aspect ratio is based on the span of the wings; however, there is some contention that since only air flow perpendicular to the quarter-chard line is considered to affect the aerodynamic characteristice, aspect ratio should be based on the length of the quarter—chord line. Such an assumption is used in the analysis included in reference 1. In figure 4, the experimental results (taken from the linear portion of the lift curvs) are shown together with the predictions based on theory for both concepts of aspect ratio<sup>3</sup>. For swept-back wings, basing the aspect ratio on the length of the quarter-chord line gives the better agreement; whereas for swept-forward wings, basing the aspect ratio on the conventional span gives the better agreement.

It is believed that neither of these aspect ratio concepts gives a correct picture of the induction effects of the vortex pattern on swept wings. It can be shown that if a wing is swept back, the induction influences of the trailing vortices on the wing should be reduced, and conversely, if a wing is swept forward, the induction influences on the wing should be increased. That is, the effective aspect ratio increases with sweepback and decreases with sweepforward for wings of constant geometric aspect ratio (b2/5).

The lift-curve slopes for the wings of this report have been estimated using the method of Fallmar (reference 5) which

It is recognized that a further aspect ratio correction, namely, Jones' edge-velouity correction should be used. The effect is small, however, compared to the errors resulting from the use of simple ewesp theory. It has been omitted, therefore, in an effort to indicate clearly the adequacy of simple ewesp theory in indicating the lift-curve slope of highly ewept wings.

Taper appears to have a strong effect on lift-curve slops due to its inherent influence on induction effects. As previously mentioned, the simplified theory strictly applies only to rectangular loading and hence the taper of the wings of the subject investigation may account for some of the discrepancy between theoretical and experimental results. In an attempt to correlate the effect of taper on the lift-curve slope of swept wings, data from previous investigations of swept wings having different taper ratios (references 1, 2, 3, 6, and 7) are shown in figure 5. For most of the investigations the wing aspect ratio (defined as b 2/8) and taper ratio did not very with sweep. For those cases where aspect ratio (b 2/8) varied with sweep, the data were corrected to the aspect ratio (b 2/8) for the unewept wing.

Examination of the data in figure 5 will reveal that, as taper ratio is decreased, the maximum value of lift-ourve slope occurs at greater angles of sweepback. The relation between taper ratio and the angle of sweep at which the maximum value of lift-curve slope occurs is shown in figure 6. The figure discloses that in order to obtain maximum lift-curve slope the taper ratio should be reduced from 1.0 as the wing is swept back and, by inference, increased from 1.0 as the wing is swept forward.

A comparison of figures 4 and 5 shows that values of lift-curve slope determined from large-scale tests show no better or poorer agreement with simple theory than values from small-scale tests. It appears that the principal disegreement between theory and experiment lies in failure of the theory to properly account for the induction effects on swept wings and that in comparison the effects of scale are relatively small.

Examination of the nonlinear portion of the lift curves and comparison with small-scale data shows that no consistent effects exist which could be attributed directly to scale effect. Those differences which do exist are small and erratic in nature and probably result from differences in plan form, wing section, and local wing roughness.

Flap effectiveness. - According to simple sweep theory, flap effectiveness decreases as cos An additional correction to account for induction effects must be applied

when comparing flap effectiveness on wings of different aspect ratio. The comments previously made regarding the effects of induction on lift-curve slope apply equally well to flap effectiveness. In fact when  $\alpha^i$ ,  $C_{L\alpha}^i$ , and  $\delta$  are measured perpendicular to the quarter-chord line, the effectiveness parameter  $\alpha^i \delta$  is unaffected by sweep and the lift increment produced by flap deflection is directly proportional to  $C_{L\alpha}^i$  or to  $C_{L\alpha}$  cosA. Hence the theoretical affect of sweep and flap lift increment may be written either in terms of sweep and aspect ratio:

$$\left(\Delta C_{L}\right)_{\hat{\Lambda}} = \left(\Delta C_{L}\right)_{\hat{\Lambda}} = 0 \quad \cos^{\lambda} \frac{\left[\Lambda/(\hat{\Lambda}+2)\right]_{\hat{\Lambda}}}{\left[\dot{\Lambda}/(\hat{\Lambda}+2)\right]_{\hat{\Lambda}=0}}$$

or in terms of sweep and lift-curve slopes:

$$(\Delta C_L)_{\Lambda} = (\Delta C_I)_{\Lambda} = 0 \cos \Delta \frac{(C_{L_{\alpha}})_{\Lambda}}{(C_{L_{\alpha}})_{\Lambda=0}}$$

where a is the angle of atteck of the root chord.

In figure 7 the experimental results are shown together with predictions made in accordance with both the foregoing relations. It can now be seen that predictions of flap lift increment made in terms of aspect ratio deviate from experiment the same as did the predictions of lift-curvs slope.

However when predictions are made in terms of lift-curvs

<sup>\*</sup>Note that in correcting for aspect ratio, the aspect ratio was based on the span. As in the case of  $C_{L_G}$ , if aspect ratios were based on the length of the quarter-chord line better agreement in flap lift informent would have been obtained for swept-back wings and poorer agreement for swept-forward wings.

slopes, the agreement with experiment is almost exact. Thus the control which  $C_{L_{\mathbf{Q}}}$  has on flap effectiveness emphasizes the importance of fully understanding the effect on  $C_{L_{\mathbf{Q}}}$  of the many factors involved; this is especially significant when flap effectiveness is considered in terms of airplans control and performance.

Since the flap lift increment is dependent upon liftcurve slops, the conclusions concerning the effect of scale on lift-curve slope apply squally well to flap lift increment. In general, it can be said that sweep introduces no new scale effect on flap lift increments measured at low angles of attack.

Meximum lift.— The effect of sweep on maximum lift of
the wing without flaps, with 0.523 span flaps deflected 60°,
and with full-span split flaps deflected 60° is shown in
figure 5. Attention is called to the fact that the wings
tested were composed largely of production wing panels with
normal roughness and irregularities such as caused by access
plates. As a result of the roughness, maximum lifts measured
on these wings may be somewhat lower than those measured on
smooth wings. However, since the measured values on the
unswept wing appear to be reasonably high for the particular
airfoil sections, it is believed that the roughness was not
sufficient to seriously reduce the maximum lift measured.

As shown in figure 5, sweep in wing plan form produces serious losses in maximum lift. However, for all but one of

the wing configurations the measured maximum lift was equal to or greater than simple theory would indicate, that is,  $C_{\rm Lmax}$  did not in general decrease proportional to  $\cos^2\Lambda$ . Furthermore, the geometric angle of attack for  $C_{\rm Lmax}$  (fig. 9) does not decrease as

which would be predicted by simple theory. It is probable that spenwise boundary-layer flow prevents stall from spreading from tip to root on the swept-back wing and from root to tip on the swept-forward wing. It is also possible that this intense boundary-layer drain allows certain sections of the wing to reach abnormally high angles of attack prior to stall.

On the unswept wing the gain in maximum lift coefficient due to flap deflection is equal to the flap lift increment at low angles of attack; whereas on the swept wings the gains in maximum lift coefficient are somewhat less than the flap lift increments realized at low angles of attack. This is particularly true for the outboard portion of full-span flaps on the swept-back wings and the inboard flaps on the swept-forward wings. (For sweepback angles greater than 30°, full-span flaps produce no greater  $C_{l_{max}}$  than do partial-span flaps.) Such decreases in flap effectiveness with sweep are disappointing but not surprising, since near stall the air flow is separated on the outboard section of swept-back wings and the inboard section of swept-forward wings.

The measured loss in  $C_{\rm L_{max}}$  due to sweep seriously limits airplane performance. With either partial or full-span flape, the loss in  $C_{\rm L_{max}}$  due to  $45^\circ$  of sweep would require increases in landing speed of approximately 20 percent.

Because of differences in taper ratio and mirfoil sections of models used for small- and large-scale tests to date, no quantitative conclusions can be drawn regording the effects of scale on Ot at various angles of sweep. In general, hawever, it can be inferred that as the angle of sween is increased the effects of scale become smaller. For instance a comparison of figure 6 of this report with figure 7 of reference 5 shows that an increase in  $O_{L_{\rm max}}$  of 0.25 is obtained at 0° of sweep when going from small- to large-scale teets. In contrast, increases of only 0.10 and 0.05 in  $O_{L_{\rm max}}$  are gained with 30° and 450 sweepback. The increase in Olympk at 0° of sweep is in general accord with what past experience has shown to be a reaconable effect of scale; whereas the increases for the swept winge fall far short of what would be anticipated from experience on straight wings. These data indicate that large-scale tests show a much more rapid decrease of  $G_{L_{\rm TIRX}}$  with sweep than do model tests. This seems true whether flaps are deflected or not. Since large-scale results tend to approach small-scale recults at large angles of eweep, considerable care should be taken in trying to setimate large-scale airplane performance from swept-wing model tests. Expectations of improving  $J_{\text{Imax}}$ commensurate with that experienced at zero sweep are not likely to be fulfilled. The importance of this problem would indicate a pressing need for swept-wing tests of a number of given models throughout the full Reynolds number range.

It should be noted that the above inferences have been drawn from results of teste of winge using conventional airfoil sections. It may well be that when sufficient data become available to make similar comparisons on winge using laminar-flow sections, the effects of roughness and Reynolds number may be markedly different.

L/D ratio. - The variation of L/D with lift coefficient is shown in figure 10 for each of the five wings with partial and with full-span flaps. That part of the drag attributed to induced drag has been corrected to the aspect ratios (b /s) of the unswept wing, that is, an aspect ratio of 4.62. The L/D values for conditions where the drag coefficient was less than 0.1 are not shown because it is believed possible inaccuracies due to lack of precise drag tare values would invalidate any conclusion drawn from euch recults. This excluded study of the most important flight epeed range for plain wings and hence the L/D values for plain wings are not shown. It is believed, however, that the resulte shown for the wings with flaps are sufficiently accurate to allow useful conclusions to be drawn as to the effect of sweep on L/D ratios for that region of the most interest, centering around gliding and landing.

The results show that at lift coefficients near a maximum lift coefficient of 0.8 the L/D for the swept wings

The use of aspect ratio based upon the span of wing rather than the length of the quarter-chord line is justified on the basis of results quoted in reforence 6.

approximates that for the unswept wing. As stall is approached the L/D ratios of the swept-back wings remain at least as high as those for the unswept wing; whereas the L/D ratios of the swept-forward wings show a rapid decrease.

### Longitudinal Characteristics

The effects of sweep on the pitching-moment characteristics of the plain wing, wing with pertial-span flaps, and with full-span flaps are shown in figures 11 to 13. The remarks which follow are based upon the data obtained on the plain wing (fig. 11) tut in general apply also to the wings with partial-or full-span flaps (figs. 12 and 13).

For lift coefficients less than 0.5, the pitching-moment coefficients vary almost linearly with lift coefficient and indicate that forward sweep moves the aerodynamic center forward (4 percent E.A.C. at 45° sweep), while sweepback moves the aerodynamic center rearward (5 percent E.A.C. at 45° sweep). At higher lift coefficients the 45° sweet wings, and to a lessor degree the -30° swept wing, exhibit an abnormal diving tendency. Similar diving tendencies of highly swept wings have been reported previously (reference 1). Such irregularities in moment characteristics do not appear scripus if considered only in terms of the elevator power available with a conventional tail. However, the effect upon statio stability and the abrupt variation of elevator position and of stick force with speed may prove objectionable to pilots. In

the case of a taillese design these irregularities would be more serious. For instance, if the 45° swept-back wing were considered a poseible design with the 0.40-span ailerons used for longitudinal control (elevons) and with neutral stability at low lift coefficients, over 30° up-elevon travel would be required to maintain trim even if the elevon effectiveness at low lift coefficients were maintained. For the 45° sweptforward wing a similar but less extreme condition existed. Smaller control angles would be required but the data indicate an abruptness of control motion which, because of the low damping in pitch, might be serious in tailless designs. In considering the longitudinal stability it should be remembered that the effects of fuselage, tip shape, slots, etc., have been disregarded. It may be, and unpublished data so indicate, that minor configuration changes will remove the diving tendency and its associated problems.

For lift coefficients just less than  $C_{l_{max}}$ , the swept wings tested, with the exception of the 30° swept-forward wing, exhibited a strong climbing moment. This characteristic is obviously undesirable since it makes inadvertent stall quite likely. In reference 8 a chart was presented which defined, on the basic of small-scale data, the boundaries of aspect ratio and sweep angle which would give a wing either a climbing or a diving tendency near stall. This chart is reproduced herein as figure 14. Also shown on this figure are the data obtained in this investigation. Based upon these data, it appears that the chart as set forth in reference 8

applies as well to large-scale as to small-scale wings; furthermore, the chart applies to swept-forward as well as swept-back wings. Insofar as the over-all shape of the pitching-moment curve is concerned large-scale tests agree generally with small-scale results with the exception of minor differences. Again it should be noted that these comparisons have been made from examination of results of investigations on wings using conventional sections. The validity of the statements regarding these comparisons is as yet unsubstantiated in cases where laminar-flow sections are involved.

### Lateral Characteristics

pinedral affect.— The variation with lift coefficient of the rolling moment due to sideslip is shown in figure 15 for the plain wings and in figure 16 for the wings with flaps. The powerful influence of sweep on the dihedral effect is immediately apparent. (A scale of effective dihedral for the unswept wing has been shown on the figures to allow convenient comparisons.) Within limits, the dihedral effect due to sweep increases in proportion to lift coefficient.

Both the 30° and 45° plain swept-back wings reached a maximum value of  $\mathcal{O}_{1\beta}$  of -0.0034 (17° effective dihedral) at lift coefficients of 1.15 and 0.85, respectively. In the case of the swept-forward wing the maximum value of  $\mathcal{O}_{1\beta}$  increased with angle of sweep, being 0.0014 for the -30° swept wing and 0.0020 for the -45° swept wing. These maximum values for the swept-forward wings occurred in both cases near a lift coefficient of 0.9. It should be noted that while the swept-back wings show much greater dihedral effect than do the swept-forward wings, this is due largely to the dihedral effect of

the unswept wing. The incremental dihedral effect is roughly of the same order of magnitude for either direction of sweep.

The maximum dihadral effect of the wing with flaps deflected is considerably higher, about 32° for the 45° swept-back wing with full-span flaps. Such extreme dihadral would make maintenance of a wings-level attitude in the landing approach almost impossible because of extrems sensitivity in roll to slight angles of sideslip. Even with adequate lateral control it is felt that a pilot would have difficulty in reacting sufficiently fast to prevent reaching excessive angles of bank.

For the case where lift is changed by changing angle of attack (flap deflection constant), simple sweep theory gives the following relation for the parameter  $\partial C_{1g}/\partial C_{L}$ :  $(\partial C_{1g}/\partial C_{L})_{\Lambda} = (\partial C_{1g}/\partial C_{L})_{\Lambda=0} - 1/4 + \frac{\tan \Lambda}{57.3}$ 

<sup>\*\*</sup>Sit is recognized that both of the terms on the right side of this equation should be modified further by a correction involving aspect ratio and edge velocity. Simple theory shows that, where asymmetrical lift exists, the corrections would be the form A/(AE+4). Again the question erises as to what the value of aspect ratio should be. Obviously the choice is more complex than simply deciding whether the span should be based on conventional span or quarter-chordline length. In attempting to correlate the subject data as well as other swept-wing data both these approaches were used. Since neither proved consistently superior to the other or to simple theory, it was decided to delete the correction entirely. It is possible that additional study of existing data togsther with future tests will reveal a means of determining an effective aspect ratio which when used in this connection will more accurately predict asymmetric loading conditions. It should be noted, then, that throughout the sections of this report dealing with asymmetric loading conditions (sideslip or allerons deflected) no corrections for aspect ratio changes have been applied to the predictions for the effect of sweep.

This relation has been used to estimate the values of  $\frac{\partial C_{i_{p}}}{\partial C_{i_{p}}}$  for the five wings tested and the results are compared with experimental values in figure 17. Reasonably good agreement is shown except in the case of the 45° swept-forward wing which had a somewhat lower value of  $\frac{\partial C_{i_{p}}}{\partial C_{i_{p}}}$  then was predicted.

For the case where lift is changed by changing flap deflection (angle of attack constant), the theoretical effect of sweep on  $\partial C_{1g}/\partial C_{L}$  is twice that given by the foregoing expression, that is,

$$\left(\frac{\partial C_{L_{\beta}}}{\partial C_{L}}\right)_{\Lambda} = \left(\frac{\partial C_{L_{\beta}}}{\partial C_{L}}\right)_{\Lambda=0}^{-\frac{1}{2}} \frac{b_{T}}{b} \frac{ten\Lambda}{57.3}$$

where b<sub>f</sub>/b is the ratio of flap span to wing span. The estimated and experimental results for this case are also shown in figure 17. The agreement between theory and experiment in this case is only fair. The disorcpancy is probably due in great measure to failure of the theory to properly account for the spanwise center of load. Theory indicates rectangular loading—that is, that the center of additional load is applied at midsomispan of wing or flap. Howement of the center of load inboard as much as 20 percent of the wing semispan would be required to make the discrepancy between theory and experiment vanish.

Thus, relations obtained by means of the simplified theory appear to estimate at least the gross effects of sweep on the parameter  $\partial C_{L_{\beta}}/\partial C_{L}$ . A notable exception is the case of the

45° swept-forward wing which exhibits a dihedral effect much less than theory would indicate but still greater than the 30° swept-forward wing.

Since the problem of determining the value of  $\partial C_{l,\beta}/\partial C_{l,\beta}$  and the maximum value of  $C_{l,\beta}$  is probably the most serious one faced by the designer of swept-wing simplenes, a considerable effort has been made to evaluate the effects of scale on swept-back wings from the data. Unfortunately, such an evaluation could not be obtained. Only the generalization can be made that the effects of scale appear much less important than the effects of wing geometry. Both large- and small-scale swept-back-wing tests show very similar characteristics. That is, the value of  $\partial C_{l,\beta}/\partial C_{l,\beta}$  (approximates that predicted by theory, with a maximum value of  $C_{l,\beta}$  being reached prior to the stall, and followed by a reduction in  $C_{l,\beta}$  has the stall is approached. Note that these and the following considerations regarding the effects of scale apply to plain swept-back wings only.

As previously noted, the value of  $\partial C_{t_{\parallel}}/\partial C_{t_{\parallel}}$  indicated by simple theory is based upon the assumption that the additional load is concentrated at the mid-semispan. Therefore marked differences in this parameter would be expected where nonrectangular loading was known to exist. In comparing experimental results with the theory such was found to be the case. Reference I showed that for the rectangular swept-back wings the measured value was as much as 14 percent more than

the predicted value. The tests reported herein gave experimental results less (as much as 14 percent less) than the predicted value. Such differences might be anticipated since theory shows that sweepback tends to shift the load center towards the tips, and taper ratios less than 1.0 tend to shift the load center towards the root. A complete understanding of this action cannot be had until more thorough studies are made of the effects of sweep and taper on load center. A first approximation (probably an overcorrection) of the answer can be reached, however, by simply adjusting the load center to correspond to the area center. If this is done, theory would fall within 10 percent of the results shown in this report while, of course, the discrepancies of reference 1 would be unchanged. Such a procedure applied to the results of reference 3 would slightly overcorrect for the effects of taper that are shown. From this it can be concluded that the value of dCi,/dCi can be approximated to within 15 percent by simple theory; that a closer approximation can be had - probably within 10 percent - if the centraid of land is assumed to lie on the centroid of area. It is believed that the effects of scale fall within this latter error and probably are of the same general magnitude as the effects of section or tip shape. Es data could be found to aid in a quantitative evaluation of these effects.

With regard to the maximum values of Cig likely to be encountered with a highly swept wing, it appears impossible to

conclude more than the fact that a maximum value exists for every wing and that this maximum value tends to decrease with taper ratio. The data of this report and reference I show very nearly the same maximum value  $\left(c_{t_{\text{finax}}} = 0.0035 \text{ to } 0.0036\right)$ for both swept-back wings. Reference 3 shows a very similar maximum for the untapered wings but shows the maximum decreasing with both taper and sweep for other wings. No relation seems to exist between the lift coefficient at which the maximum Ct, occurs and the maximum lift coefficient of the wing. Since, however, the value of dCLE/dCL, in general, increases more rapidly with sweep than CLmax decreases, the maximum value of Cig occurs at progressively lower percentages of  $C_{\mathrm{Lmax}}$  as sweep is increased. For instance, for a 45° sweptback wing,  $c_{l_{max}}$  occurs at 0.55  $c_{l_{max}}$  in reference 3, 0.61  $c_{l_{max}}$  in reference 1, and 0.70  $c_{l_{max}}$  in the data of this report; whereas Ct fnax for a 30° swept-back ring occurs at 0.80 CLmax, 0.82 CLmax, and 0.91 Clmax for the same data, respectively. Since the phenomenon which causes the valve of Ct, to peak are not completely understood at the present time, an accurate prediction of its value is impossible. Examination of all available data leads, however, to the conclusion that if a maximum value of -0.0033 is chosen the choice can be considered conservative, but for the present, wind-tunnel tests must be relied upon to give the exact answer. Certainly this problem is worthy of additional study. Until the governing factors are more clearly defined it remains impossible to determine to what extent  $c_{l_{\beta_{\max}}}$  is affected by scale.

In the case of swept-back wings with flaps deflected the amount of correlation possible between large- and small-scale data is extremely limited and the results far less amenable to interpretation. For most cases examined theory gave at least a slightly conservative value of dClg/dCL where the change in lift coefficient was due either to a flap deflection at constant angle of attack or a change of angle of attack with flaps deflected. All the data, large and small scale showed or indicated that a value of Clamax existed and that it increased with sweep. While no systematic variation of Cipmax with wing geometry could be ascertained, none of the data showed a value greater than -0.007. For the present, therefore, if wind-tunnel tests are not available the best approach to predicting Cle characteristics of swept wings with flaps is to use simple theory to predict 6018/00L and to consider -0.007 Cls as the maximum.

No correlation was attempted with the swept-forward-wing data because of the secreity of low-ecale tests.

Alleron effectiveness. The varieties of mileron effectiveness with lift coefficient is shown in figure 15. The values of mileron effectiveness shown in figure 15 were obtained as the  $\Delta C_1$  produced by  $-15^\circ$  or  $15^\circ$  of mileron deflection and hence are  $\Delta C_1/\Delta \delta_a$  rather than a true  $C_1\delta_a$ . It is immediately apparent that mileron effectiveness decreases with sweep,

decreasing as much as 50 percent for 45° of sweep. The effectiveness of the allerons on the swept-back wings decreased with lift coefficient, rapidly at high lift coefficients.

This is due to a loss in effectiveness of the upward deflocted alleron which is in the wake of the separated flow and hence contributes little or nothing to the rolling moment. The allerons on the swept-forward wings show a general increase in effectiveness with lift coefficient, probably due to a favorable effect of the spanwise boundary-layer drain.

According to the simplified theory, as a wing is swept, the aileron effectiveness will decrease, as for any flap, in proportion to the  $\cos^2 A$ . That is, when corrections for aspect ratio are ignored the value of  $Ct_{\delta_A}$  is given by the following relation: (See footnote 6, p. 21)

$$\left( {^{\text{C}}}{i}_{\delta_{\mathbf{a}}} \right)_{\Lambda} = \left( {^{\text{C}}}{i}_{\delta_{\mathbf{a}}} \right)_{\Lambda} = 0$$

This relation has been used to predict the variation of aileron, effectiveness with sweep for the five wings tested, and the results are compared in figure 19 with the experimental data cross-plotted from figure 18 for zero lift.

The allerons on the wings of this investigation varied both in the relative amount of wing area affected and in the relative spanwise location of the center of pressure of the area affected. In comparing theory and experiment, these variations were accounted for by correcting the theoretical values of alleron effectiveness in proportion to the ratio of the relative area and spanwise center of pressure of the swept wing to the relative area and spanwise center of pressure of the unswept-wing.

For either sweepforward or sweepback the experimental values of alleron effectiveness are as much as 20 percent lower than the theoretical values.

The foregoing results show that aileron effectiveness is reduced by wing sweep and that on swept-back wings the aileron effectiveness is further reduced at high lift coefficients. Insofer as rolling control at low lift coefficient is conversed, theory shows that  $C_{\mathbf{l}_{\mathcal{D}}}$  is reduced in proportion to cosA; whereas Cloa is reduced in proportion to cos A and hence pb/2V will be reduced in proportion to cosA for a given size of aileron. In general, therefore, it appears that, to maintein a given value of pb/2V, aileron size must be increased as wings are swept. As higher lift coefficients are reached the lateral-control problem becomes particularly pronounced. Not only must powerful lateral control be provided to overcome great dihedral effects but the results reported herein show that available lateral control, at least for swept-back wings, decreases seriously with lift coefficients. For example, with the 45° swept-back wing equipped with fullspan flaps and flying near Clmax, 13° of total aileron deflection would be required to hold the wings level for only 1° of sidsslip. The need for development of adequate aileron control or a means to reduce elleron control requirements is , sucivác

# Directional Characteristics

The variation with lift coefficient of the yaving noment due to sideslip is shown in figure 20 for the plain wings and in figure 21 for the wings with partiel- and full-span flaps. Sweepback increased the directional stability and ewecoforward decreased the stability; however, due to the initial positive stability of the unswept wing the stability of the sweptforward wings became negative only at higher lift coefficients and then only slightly so.

The theoretical effect of sweep on the directional stability is in accordance with the following relation:

The directional stability estimated on the basis of this equation is compared with experimental results in figure 22. Although precise agreement is not obtained, the trend of the experimental data is indicated by theory.

The directional charactoristics of the swept wings tested should not present any serious problems of a purely low-speed static directional stability and control nature since adequate stability and control should be obtainable by use of fins and rudders of normal proportions combined with normal tail lengths. However a dynamic problem arises from the fact that  $c_{1_{ar{p}}}$  increases with sweep more rapidly than  $c_{n_{ar{p}}}.$  This unbalance between Cig and Cas leads to the dutch-roll type

of instability which has been discussed in reference 8. The data obtained in the investigation reported herein substantiate previous tests conducted on small-scale models and consequently indicats that means must be found to balance  $C_{1\beta}$  and  $C_{ng}$ .

### CONCLUDING REHARKS

Large-scale tests indicate that the primary problems to be overcome before successful use can be made of high angles of sweep are (1) high dihedral effects accompanied by poor lateral control at high lift coefficients, (2) low maximum lift value together with low flap effectiveness, and (3) rapid shift in neutral point in the moderate to high lift-coefficient range coupled with a possibility of strong stalling moment at maximum lift resulting from poor plan-form choice.

In general, simple theory enables good predictions to be made of the gross effects of sweep on wing characteristics, but it is felt that the accuracy is inadequate for purposes of design. It appears that the majority of the inaccuracies result from an incomplete understanding of the effects of aspect ratio.

Mhere it has been found possible to compare Jarge-scale data with small-scale data a comparison has shown that where scale effects exist at low angles of sweep, scale effects

tend to vanish at high engles of sweep with large -scale results approaching small-scale results.

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> Gerald M. McCormack, Mechanical Engineer.

Victor - Stevens, Jr. Victor I. Stevens, Jr. Aeroneutical Engineer.

Approved:

John F. Fersons, Aeronautical Engineer.

#### APPENDIX

Description of Basic Wind-Tunnel Test Results

For each swept wing, six-component force data were obtained at several angles of sideslip and several values of dynamic pressure. (See fig. 23 for variation of Reynolds number with dynamic pressure.) At each angle of sideslip, several model configurations were tested including plain wing, wing with partial-span split flaps, wing with full-span split flaps and wing with split-flap-type alleron. The data obtained are presented in figures 24 to 91 in terms of the variation of the measured characteristics with lift coefficient. Table II forms an index of these figures presenting the basic data.

All the data are referred to the stability axes whose origin is located at a point on the root chord or root chord projected and at the same fore and aft location as the quarter N.A.C. The test results are presented in the form of standard NACA coefficients as defined in the section Coefficients and Symbols.

All the basic wind-tunnel data have been corrected for air stream inclination and for wind-tunnel-wall effects. A brief analysis of the effect of sweep on tunnel-wall corrections indicated that the average correction either with or without sweep was approximately the same for the tunnel wing configurations considered. Hence the standard corrections for unswept wings were applied.

Force tests made with the sting support clone in the tunnel showed that its tare should be negligible except in the case of pitching moment, drag and yawing moment. Heasured pitching-moment tares are believed reliable and were applied to all the data. While the drag teres are appreciable (approximately 0.02 in the case of the unswept wing where the area is small and decreasing for the swept wings where the area is larger), it is felt that they could not be determined with sufficient accuracy to werrent application. Hence no drag teres have been applied. Since the mersured yawing-moment tares (fig. 92) were small, they were not applied to the basic data. However, in analysis of the data it was found that the tares were relatively large then compared to the effects of sweep. In order then to properly assess the effects of sweep, it was necessary to apply tores to the summary data which is, therefore, shown fully corrected in figures 20, 21, and 22.

11 .

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TABLE 1.- GROUPERIC CHARACTERISTICS OF THE FIVE STEPT MINGS

Spen of miler- ons (% b)	34.7	30.9	30.2	7.45	55.9
Spen of full-spen flaps (% b)	97.0	93.2	92.6	97.0	96.2
Span of partial- spen flaps (% b)	62.3	62.3	62.5	62.5	62.3
illeen aero-2 dynamic chord, c	11,52	3.44	6.92	7.97	10.00
Taper ratio,	0,38	3.	- 55	¥.	.42
Aspect ratio, b3/S	3,12	4.69	4.62	4.84	3.64
Area, S (sq ft)	336.5	232.3	201.8	269.4	309.6
Spen (rt)	52.58	36.39	30.53	36.06	32.56
Angle of sweep (deg)	-45	97	0	08	45

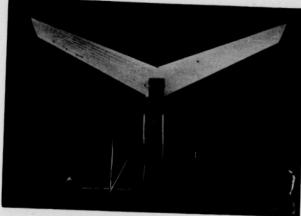
1All chords were measured parallel to the mir stream.

27laps and allerons were 20 percent chord.

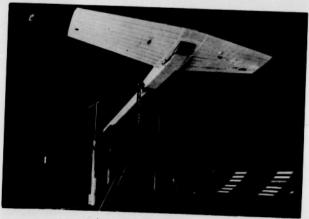
TABLE II.- INDEX TO THE BASIC DATA FIGURES

Mein wing + partiel-span fleps Co.a.cm. Co.a.ch. or vs C. To sa Apeupe to tue Co To sa Anethelled or vell CDodoCnoclocacy vs CL Plain wing + partial-span flaps | Cp. C. Cm. Cl. Cp. Cy vs CL To sa Ageugelgengeneug Plain wing + partial-span flaps | Cp. C. Cm. Cl. Cn. of vs C. CD. C. Cm. Cl. Cn. Cr. ve Cr. Plain wing + partiel-span fleps | Co.a.cm.cl.scn.cr vs CL Ched. Caschedage vac Check, Cast of De Town Va CL Date presented To sa toemorn To sa loemoto To sa 10 " D' Plain wing + full-spen flaps Plain wing + full-spen flaps Plain wing + full-spun flaps Configuration Plain wing + ailerons Plain wing + ailerons Plain wing + ailerons Metual angle of aidealip 's noted on each ourve sheet, Plain wing Plain wing Plain wing Plain wing sideslip "Nominal co 2 -10 -45 A--30 A- 0 A-30 A-45 16 Pigure number 1 63 49 48

Fig. 1a,b



(a) The 45° swept-forward wing.



(b) The unswept wing.

Figure 1.- Photographs of three of the ewept wings mounted in the Ames 40- by 50-foot wind tunnel.

Fig. 1

(c) The 45° swept-back wing.

Figure 1 .- Concluded.

Fig. 1

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Figure 1.- Concluded. (c) The 45° swept-back wing.

Fig. 2a

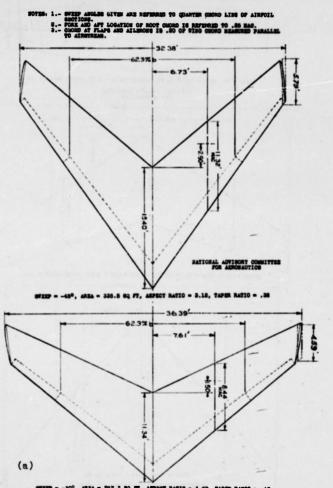
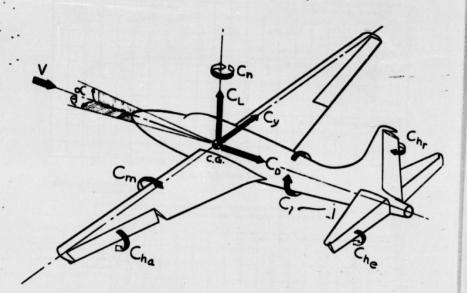


Figure 2a,b.- Geometric characteristics of the swept wings.

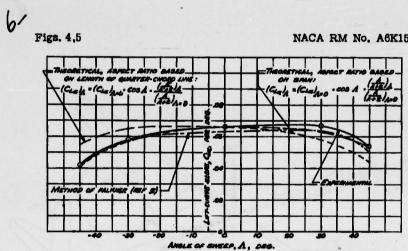
Fig. 3



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Figure 3.- Sign convention for the standard NACA coefficients.

All forces, moments, angles, and control surface
deflections are shown as positive.



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FIGURE 4.—COMMEISON BETTHEEN THEORETICAL AND EXPERIMENTAL EFFECTS OF SHEEP ON LIFT-CHAIR SLOPE

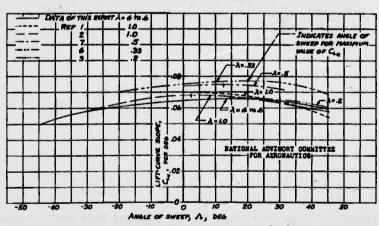


FIGURE 5. - EFFECT OF SWEEP ON LIFT-CURVE SLOPE

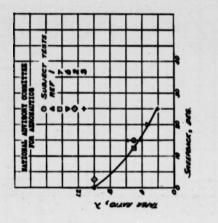
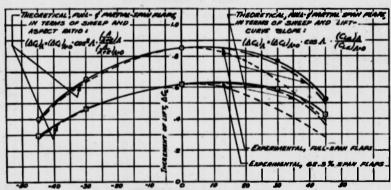


FIGURE 6- COMET SUMMERING THE BEFORE OF SMELL AS WHICH THE MACHINE MACHINE MACHINE ACCURE SLAPE SMELLING.



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FIGURE T. - COMMERSON BETWEEN THE THEORETICAL AND EXPERIMENTAL EFFECTS OF SWEEP ON THE MOREMENT OF LIFT DETAMBLE AT OF AMOLE OF ATTACK BY DEFLECTING FULL-SAM AND

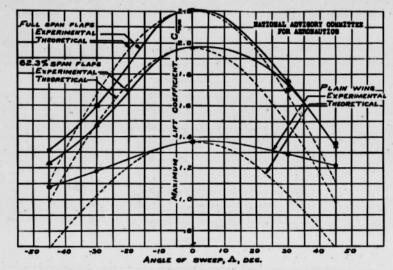
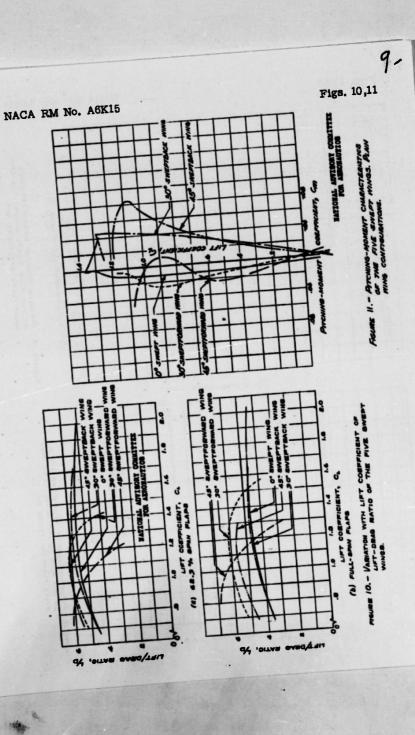
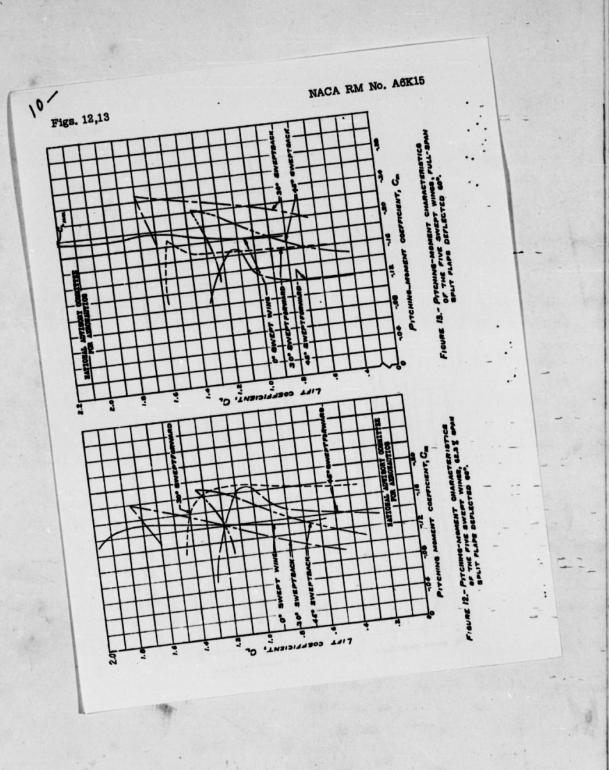
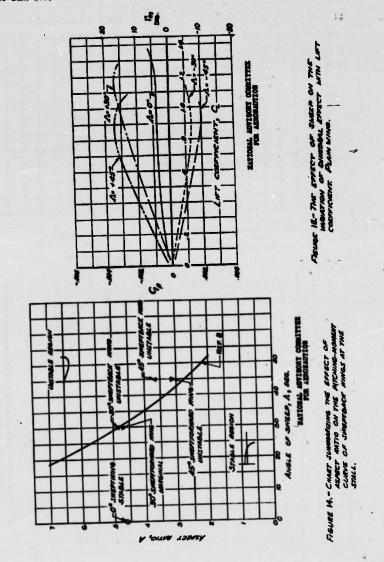


FIGURE 9.— EFFECT OF SWIFF ON ANGLE OF ATTACK FORES MAXIMUM LIFT AND FOR MAXIMUM LIFT.

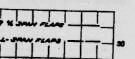


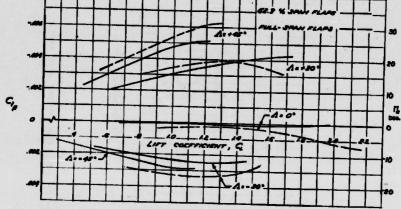


Figs. 14,15



Figs. 16,17

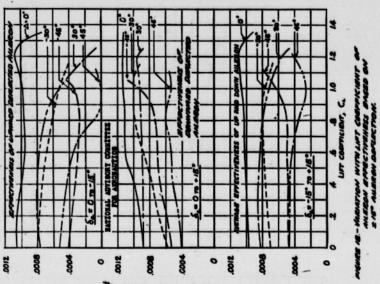




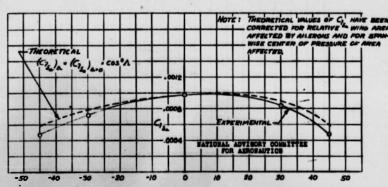
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Figs. 18,19

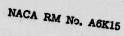


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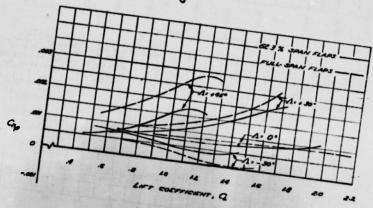


ANGLE OF SWEER A, DEG

FIGURE 19.- EFFECT. OF SWEEP ON THE ALERON EFFECTIVENESS FOR THE PLAIN WING AT ZERO LIFT,

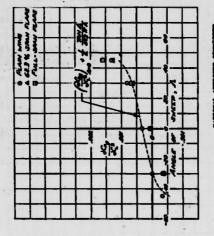






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Figs. 22,23



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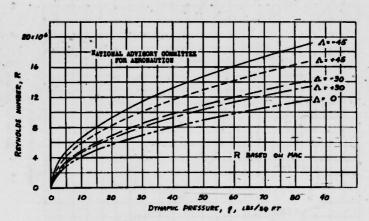
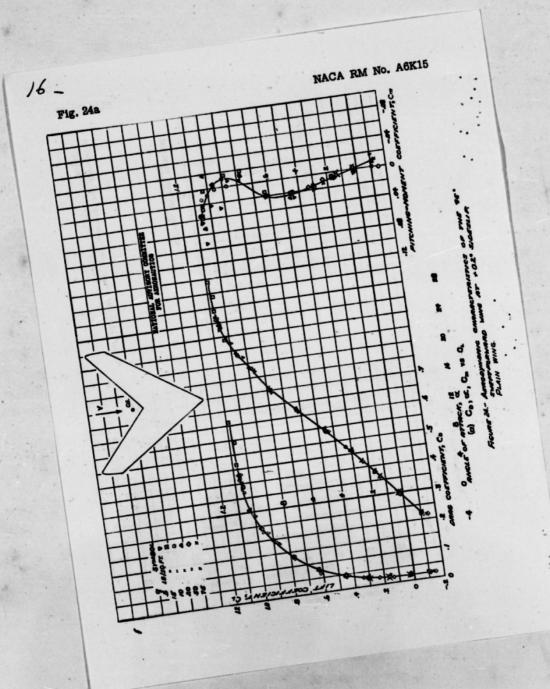
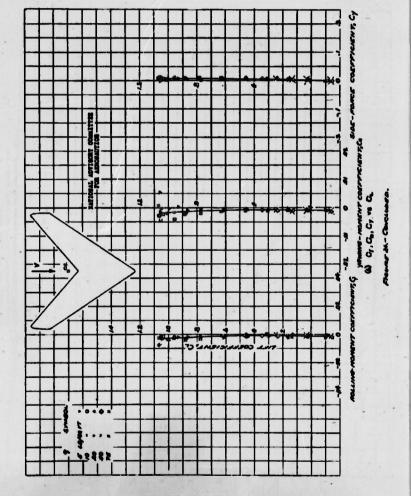


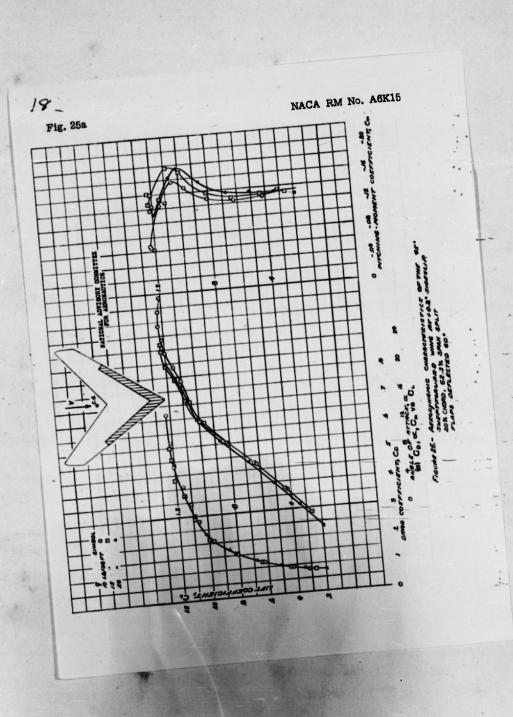
Figure 28. - VARIATION OF REYNOLDS HUMBER WITH DYNAMIC PRESSURE FOR WINGS TESTED



17 -Fig. 24b

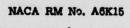
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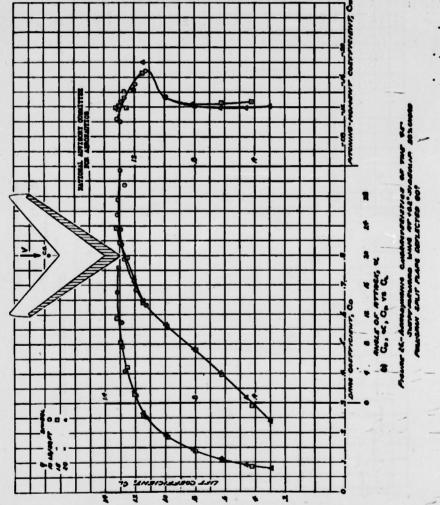




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Fig. 26a





21-

Fig. 26b

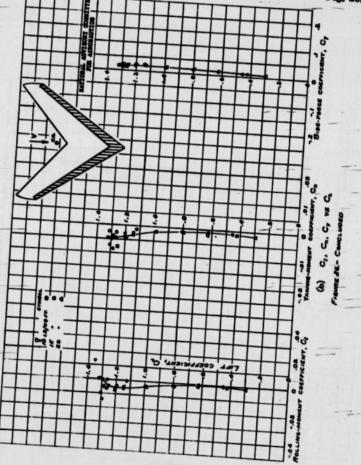
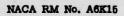
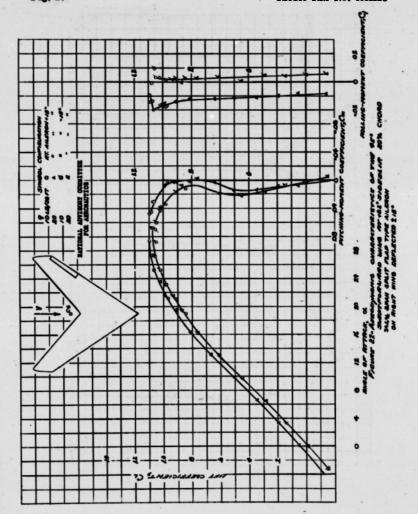
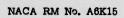


Fig. 27









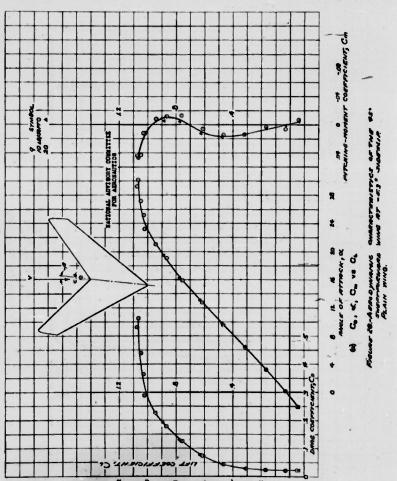
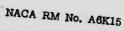
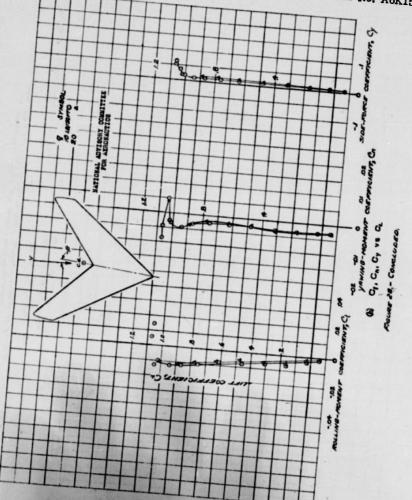
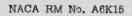


Fig. 28b

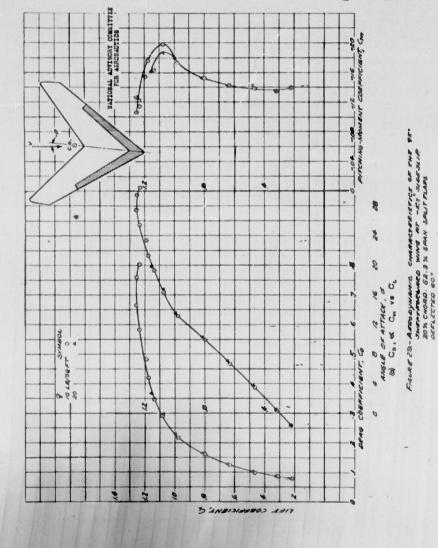




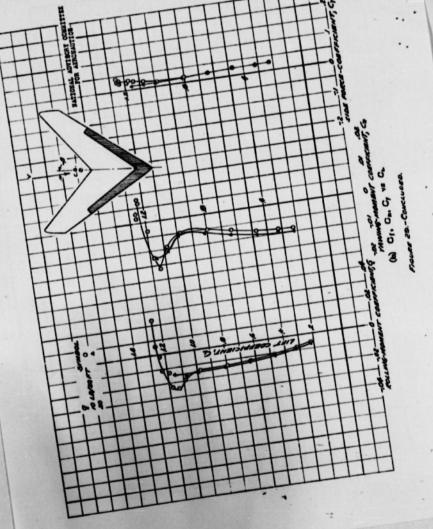


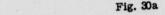






26 -Fig. 29b NACA RM No. A6K15





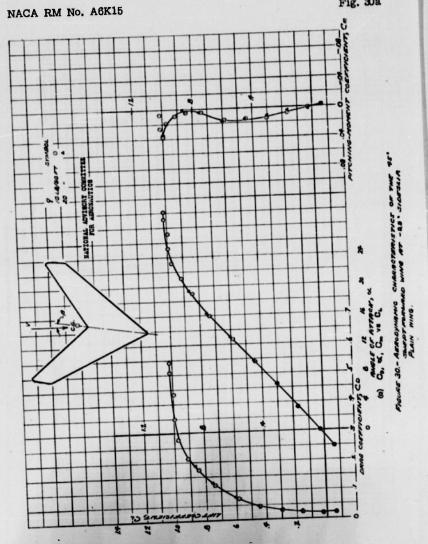
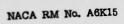
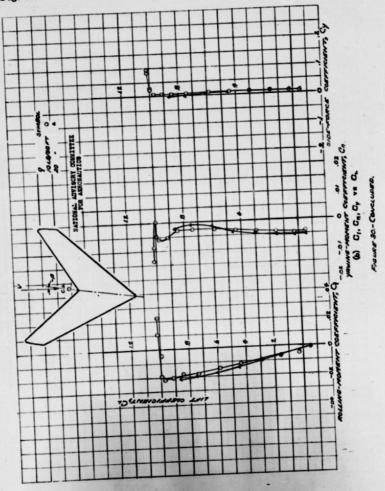
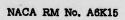


Fig. 30b









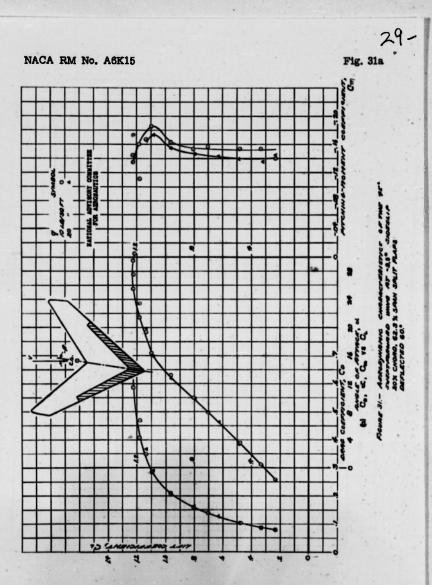
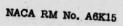
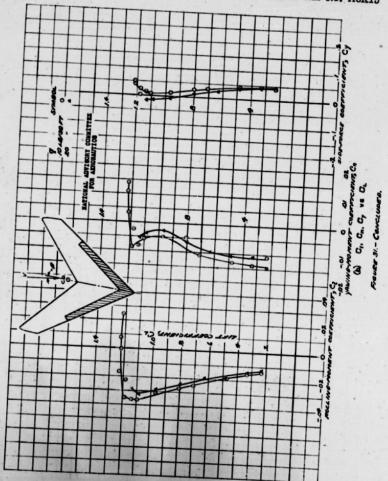
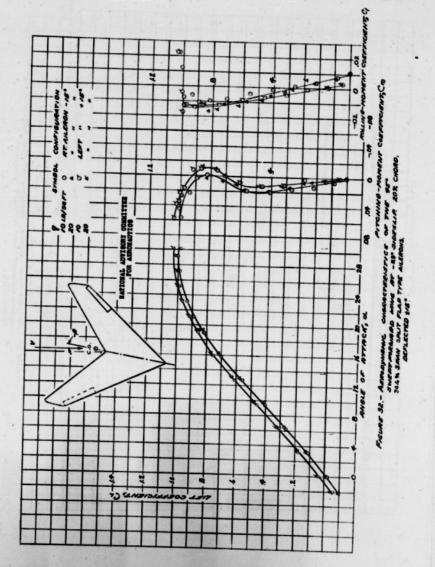


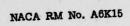
Fig. 31b

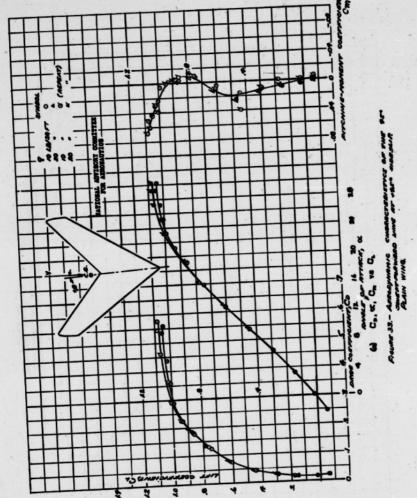






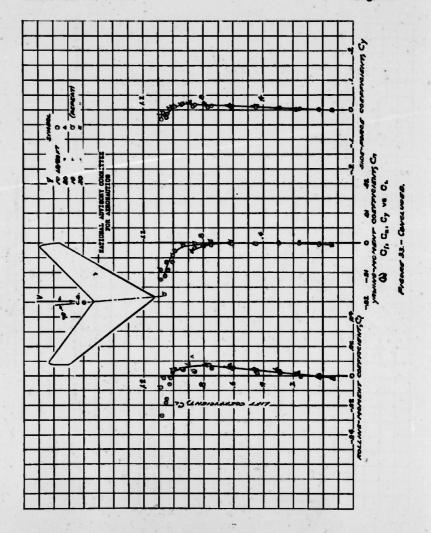






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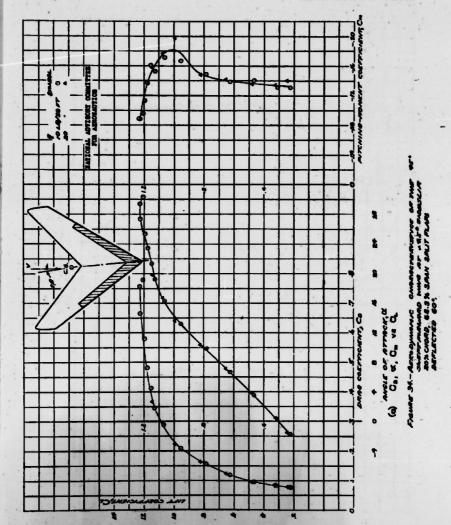


Fig. 34b

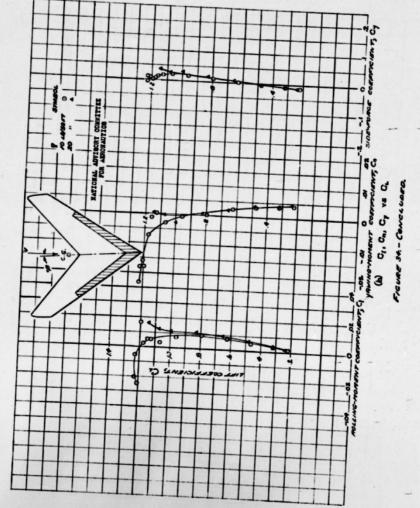
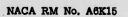


Fig. 35



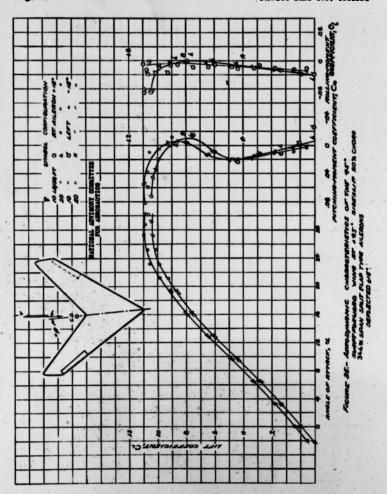


Fig. 36a

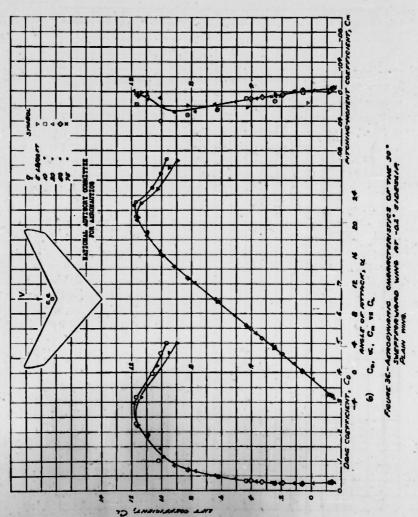
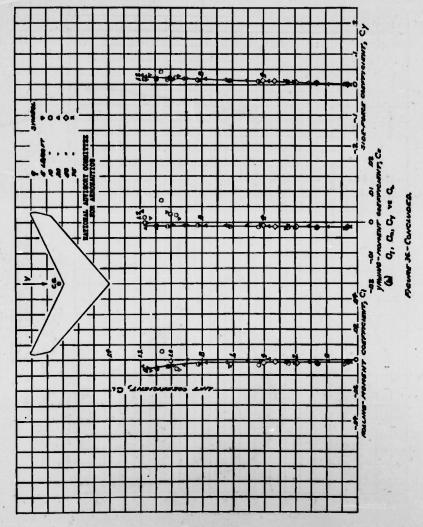
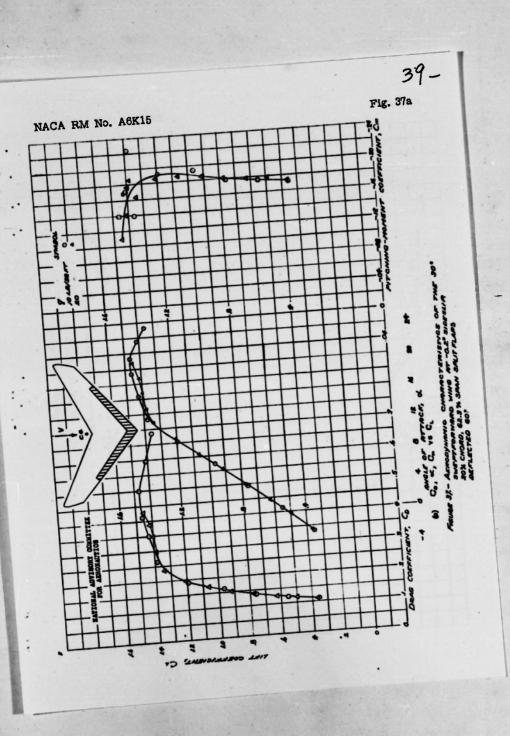


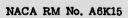
Fig. 36b

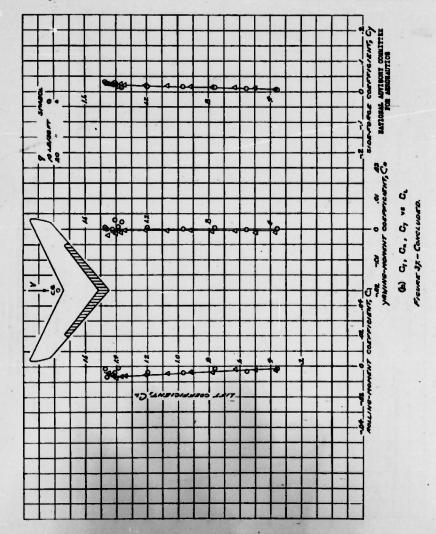
NACA RM No. A6K15





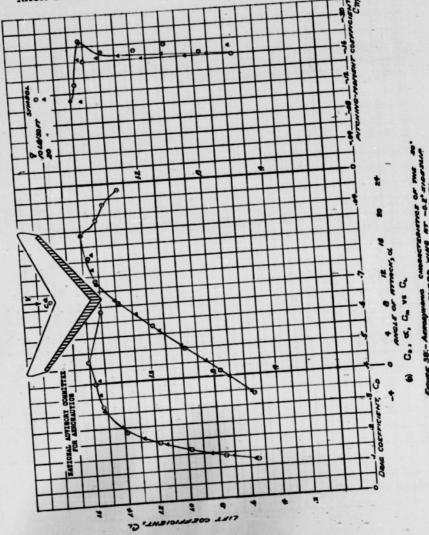
40\_ Fig. 37b



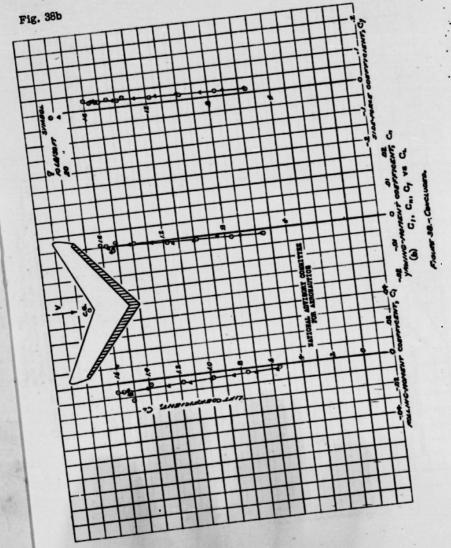


41-

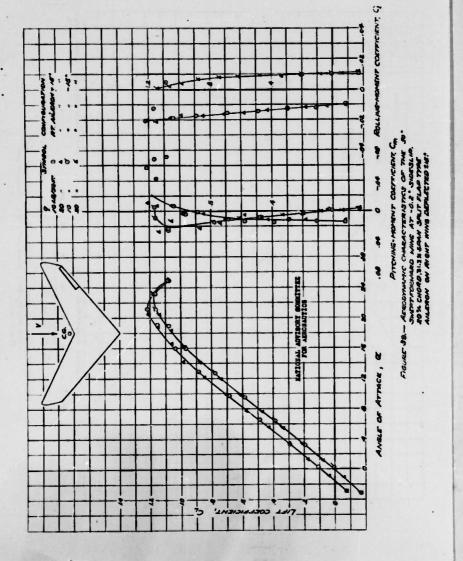
Fig. 38a



NACA RM No. A6K15



43-Fig. 39



44\_ Fig. 40a

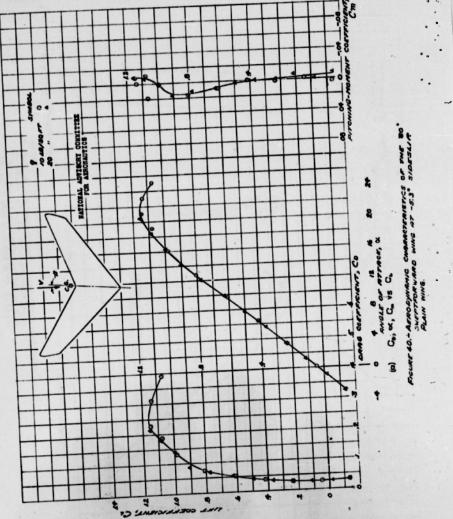
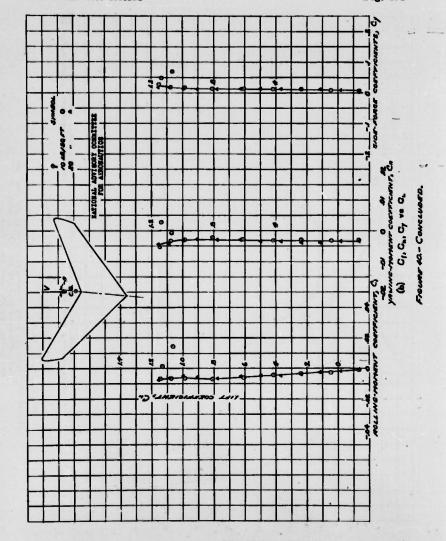
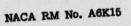
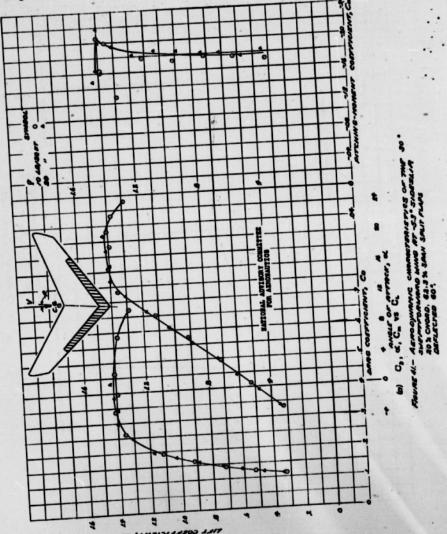
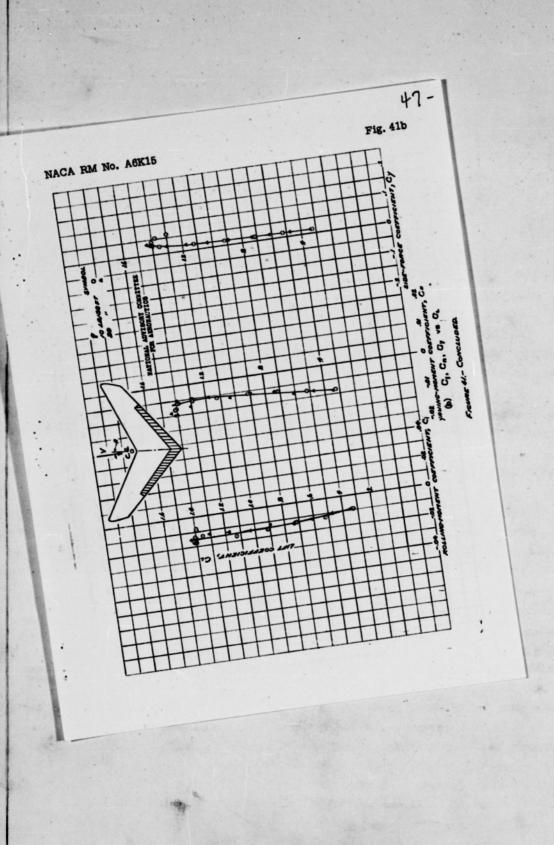


Fig. 40b



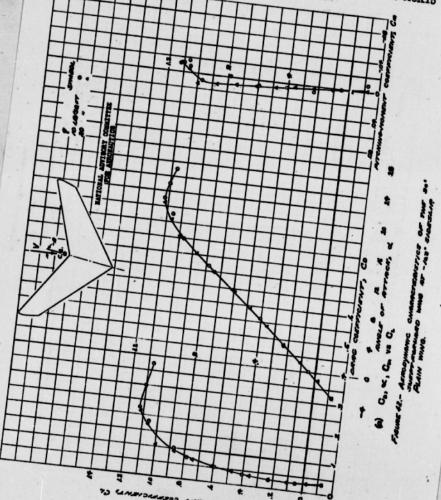






48\_

Fig. 42a



NACA RM No. A6K15

Pig. 42b

Samuel Control of the control of the

Fig. 43b

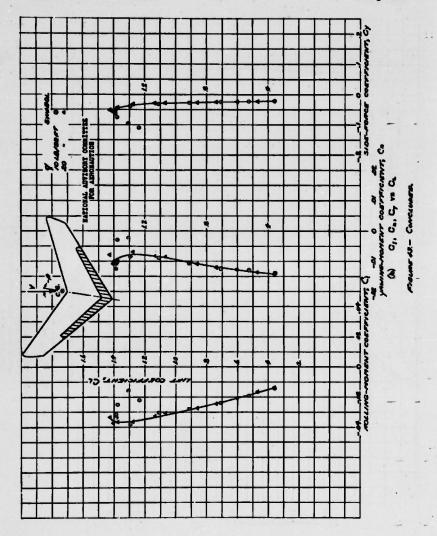
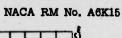


Fig. 44a



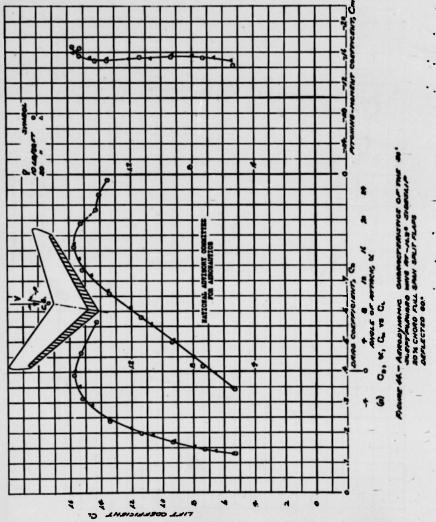
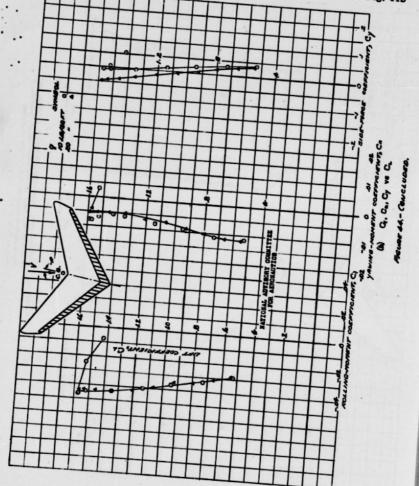


Fig. 44b



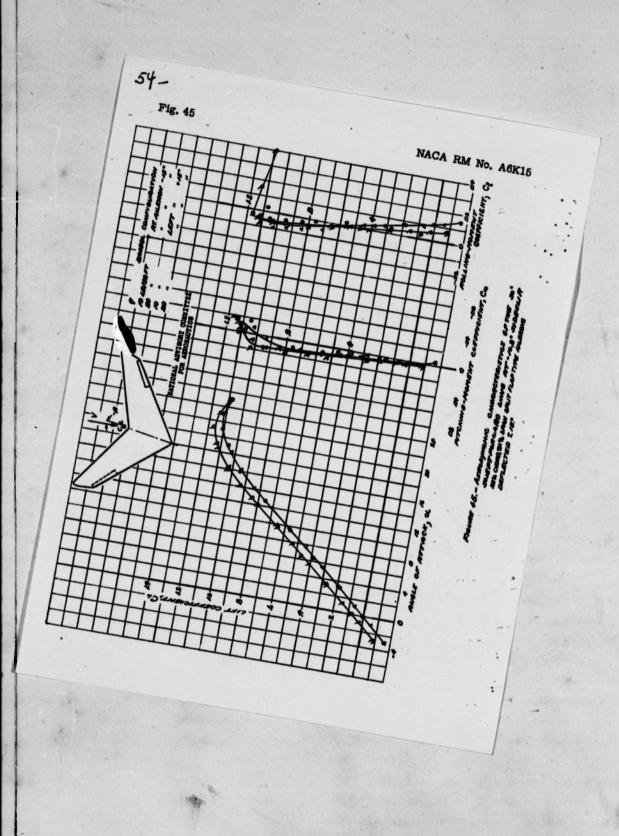
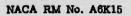


Fig. 46b



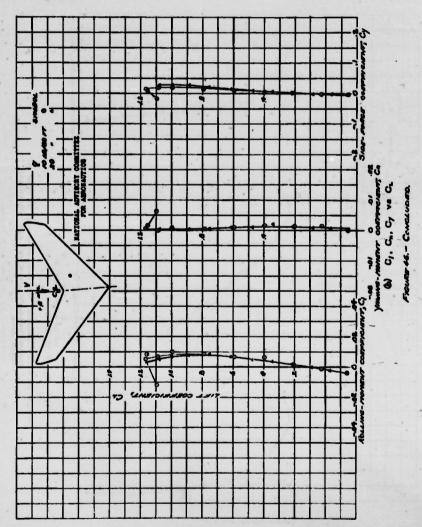


Fig. 47a

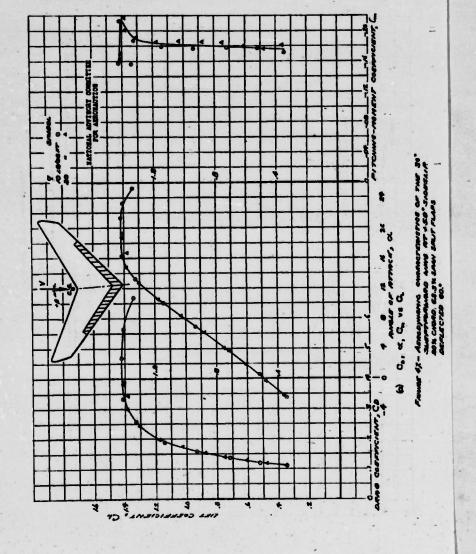
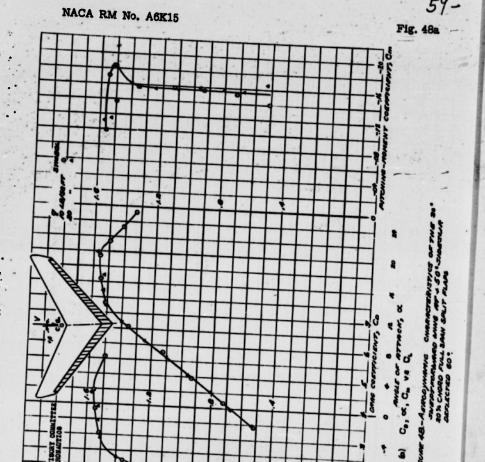


Fig. 47b

NACA RM No. A6K15

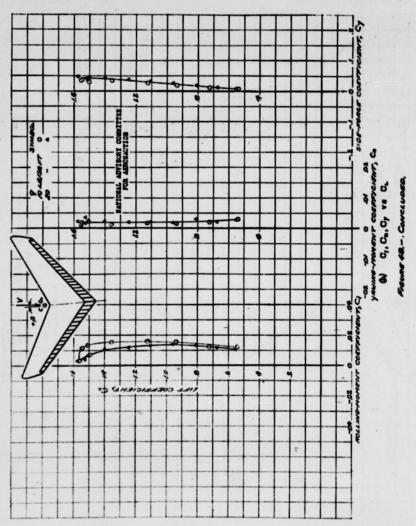
Fig. 47b

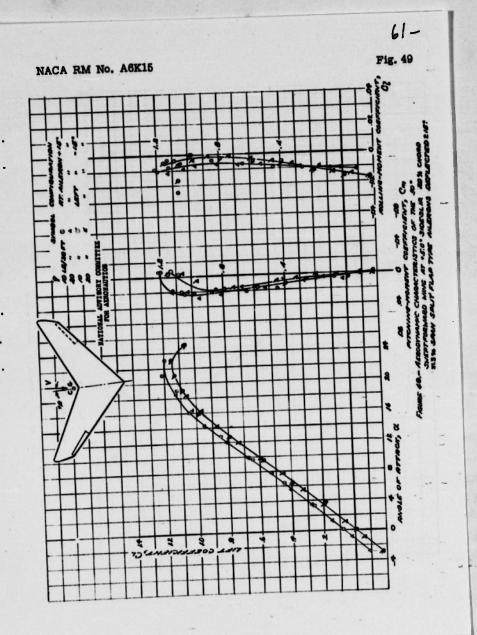


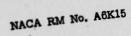
יונגו בספת מוכופ אוני כד

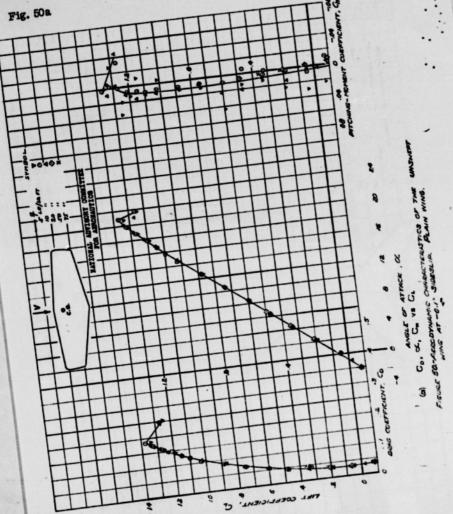
60-

Fig. 48b



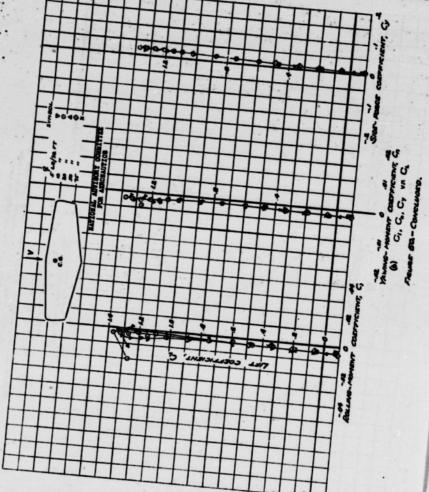






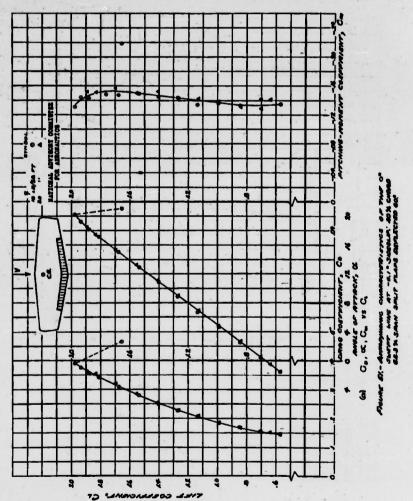
63-

Fig. 50b



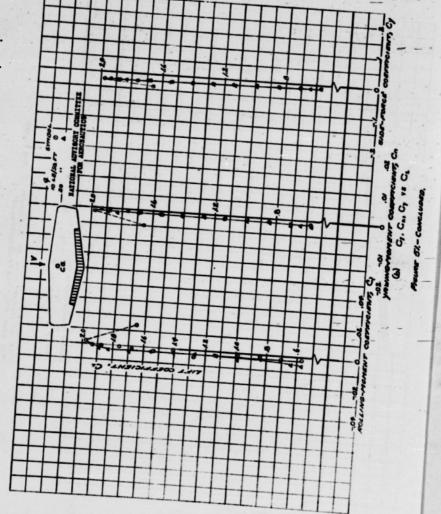
64

Fig. 51a



NACA RM No. A6K15

Fig. 51b



NACA RM No. A6K15

Fig. 52a

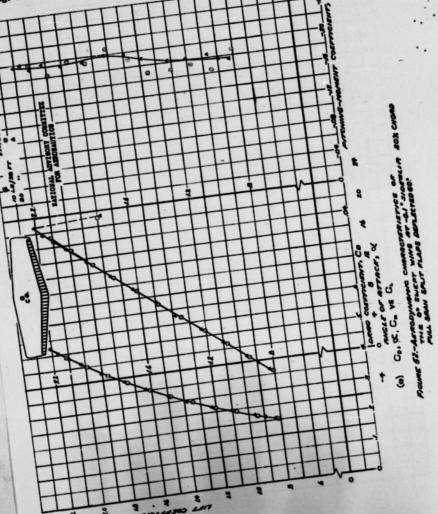
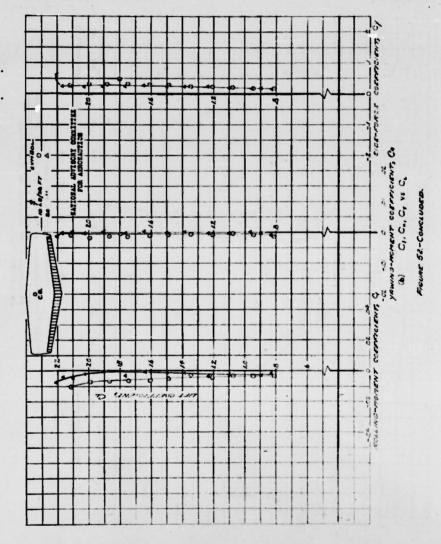
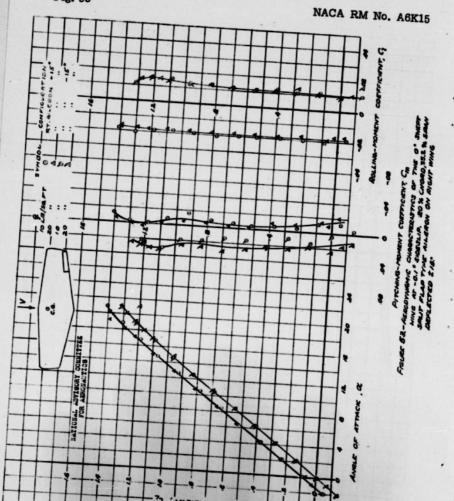


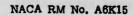
Fig. 52b



68\_

Fig. 53







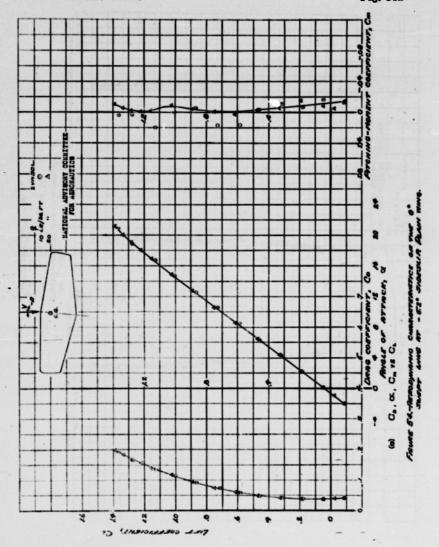
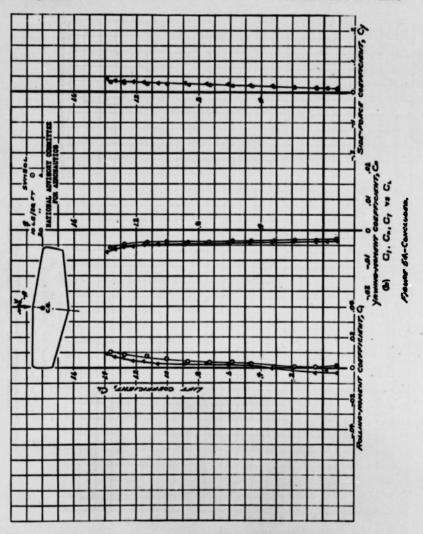
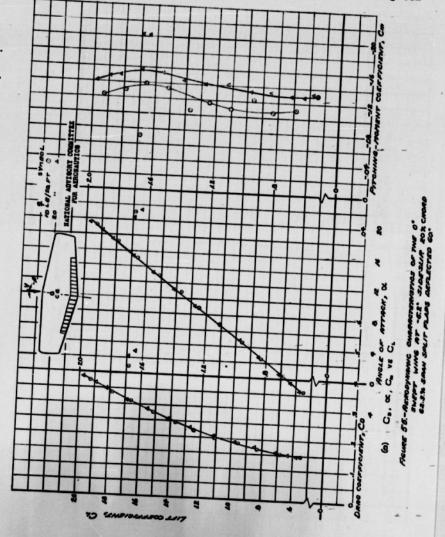


Fig. 54b







72 -Fig. 55b NACA RM No. A6K15

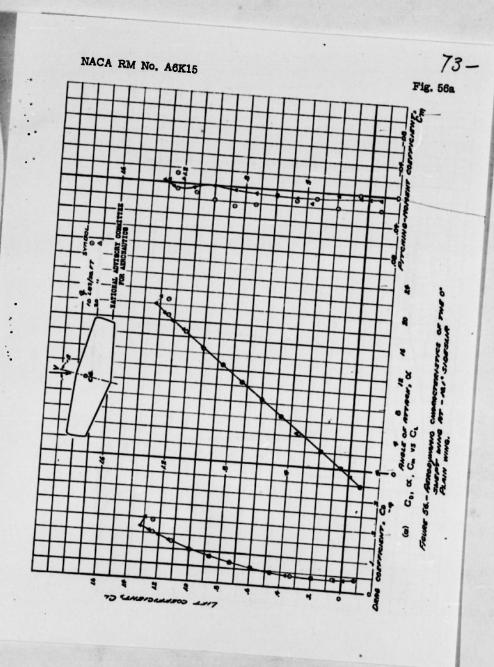
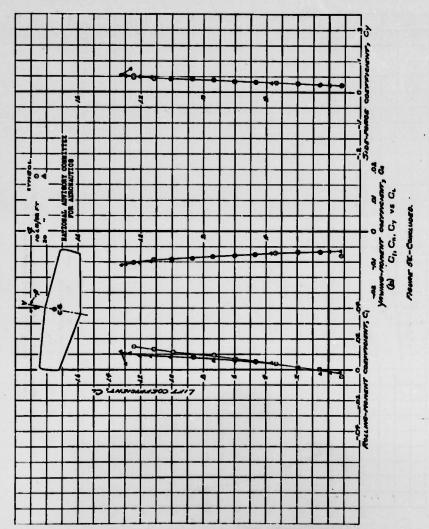
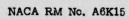
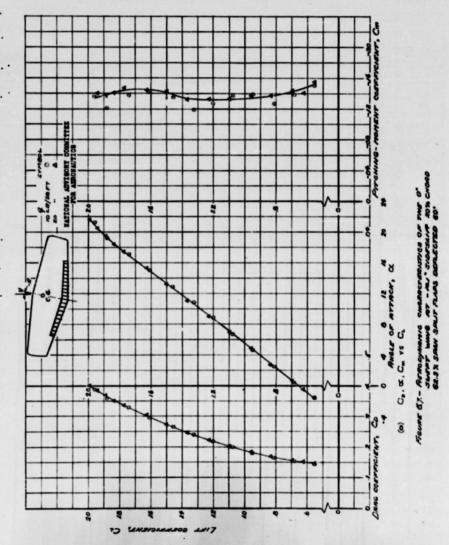


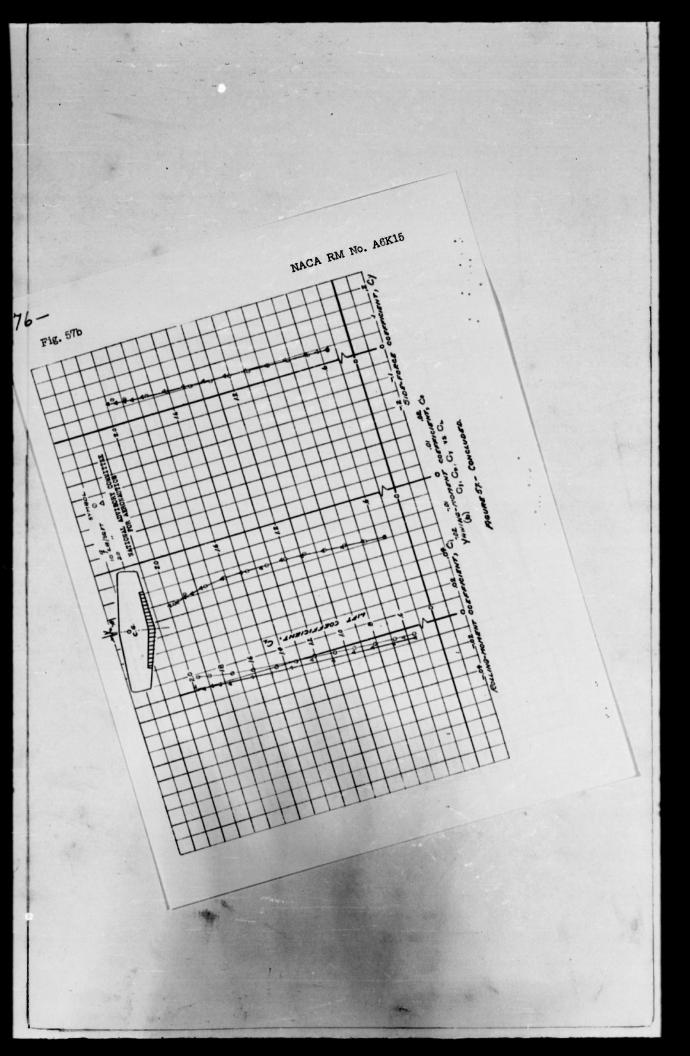
Fig. 56b





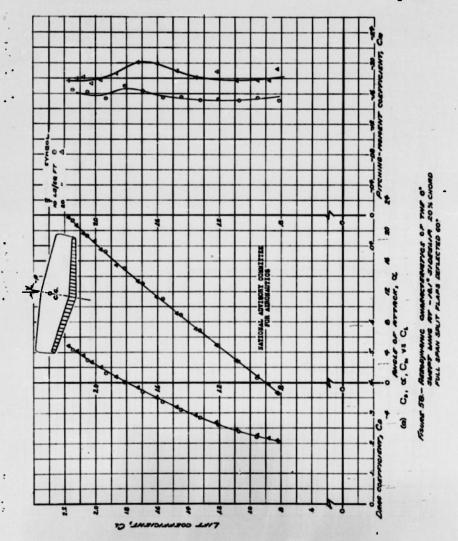


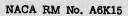




NACA RM No. 'A6K15







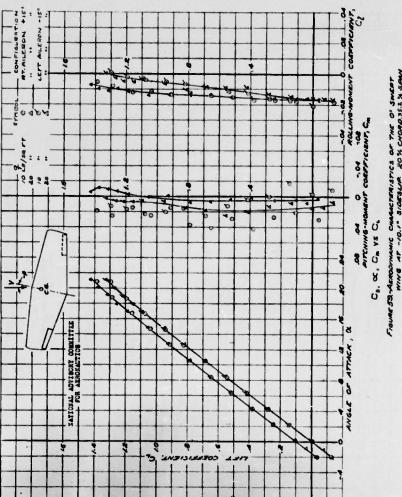


FIGURE 53-AGRODINAMIC CHARACTERISTICS OF THE O'SWEPT WING AT -10.1° SIDESLIP. 20% CHORAJSELY GAN SPLIT FLAD TYPE ALERONS DEFLECTED 2.15:

Fig. 60a

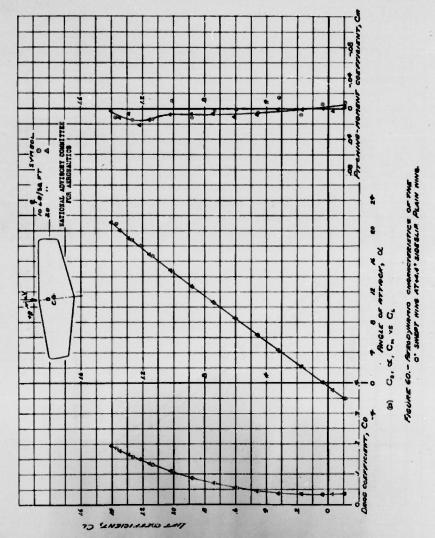


Fig. 60b

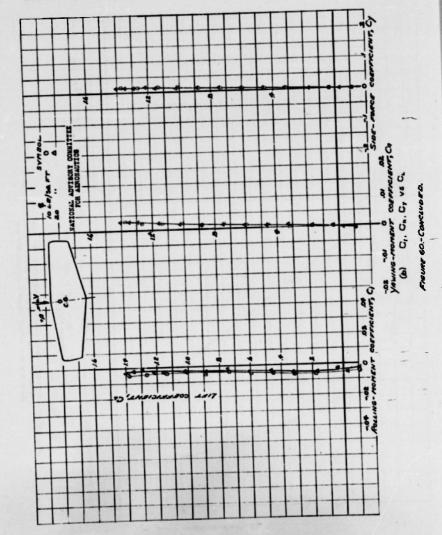
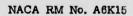


Fig. 61a



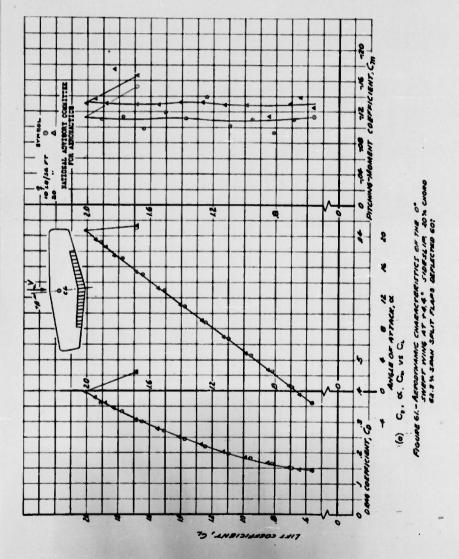
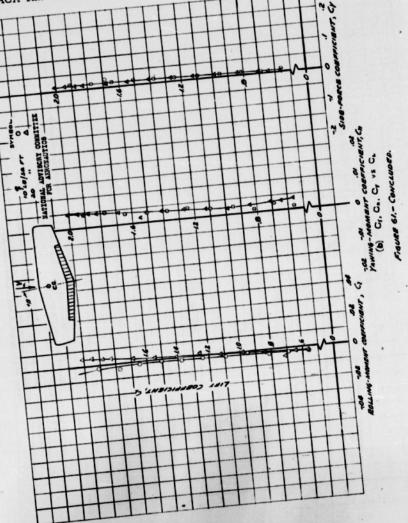


Fig. 61b



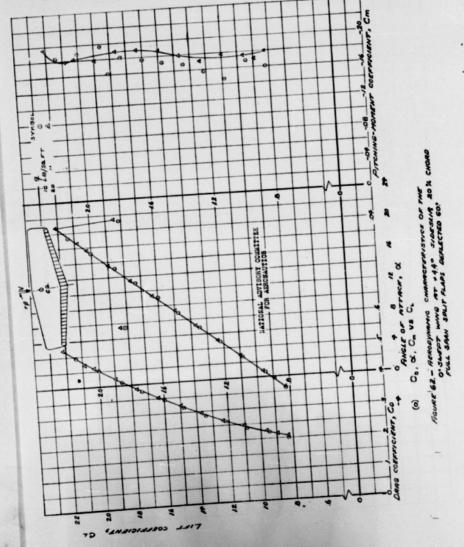


Fig. 62b

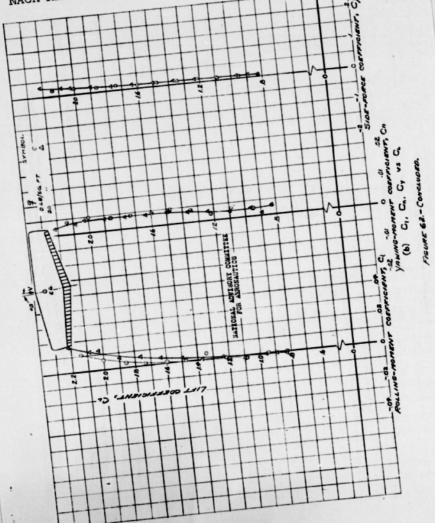
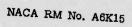
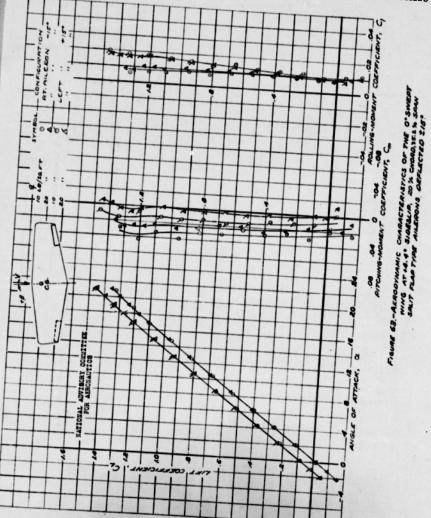


Fig. 63







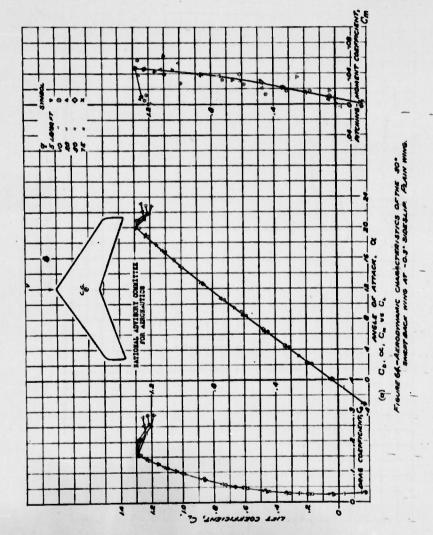
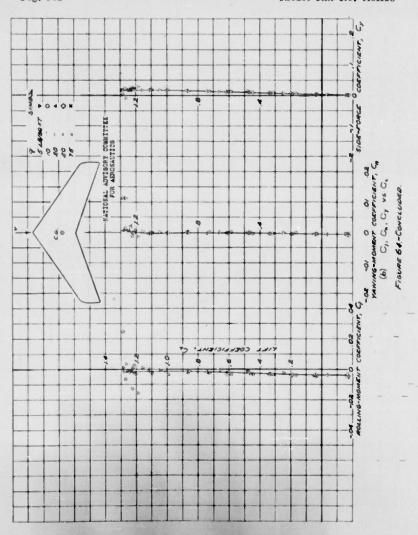
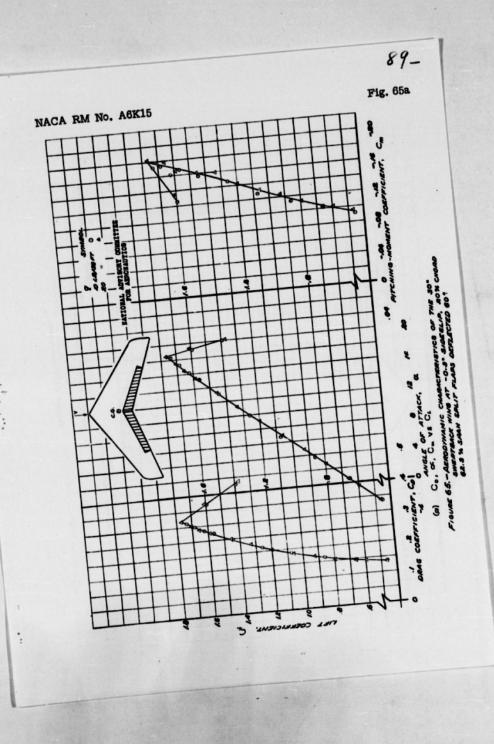


Fig. 64b

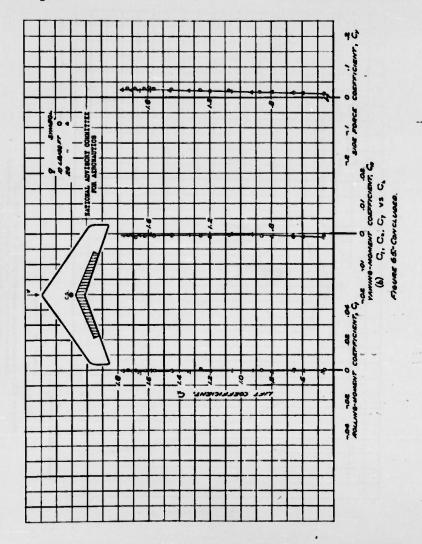
NACA RM No. A6K15





90/

Fig. 65b



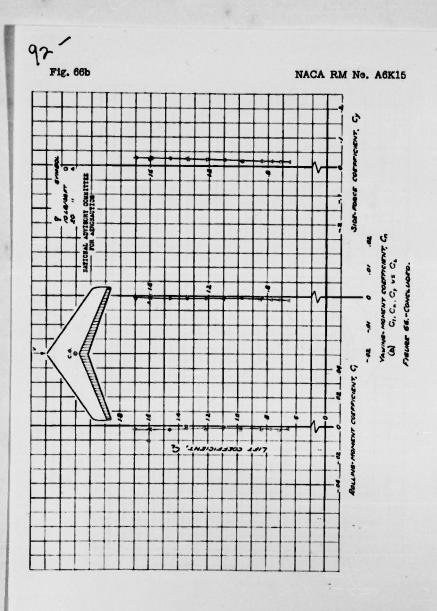
3

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NACA RM No. A6K15

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0,10



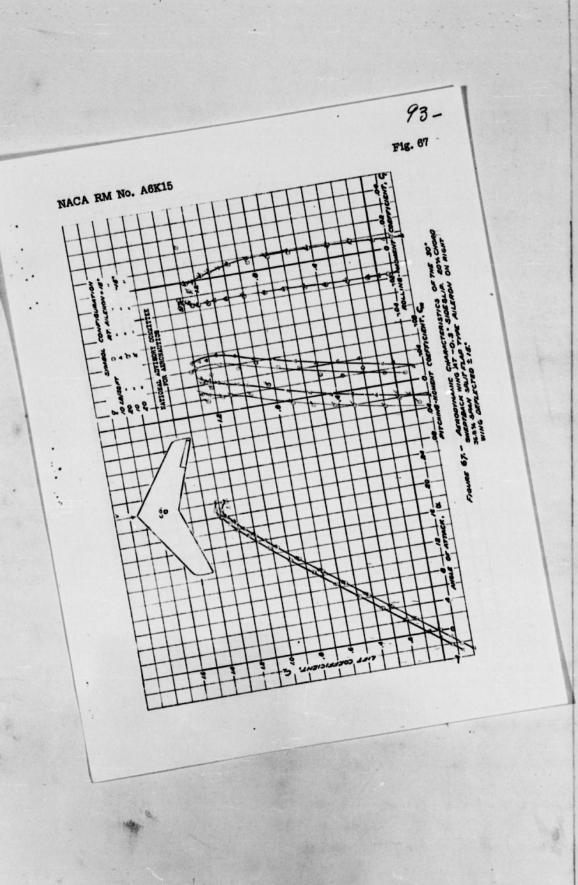
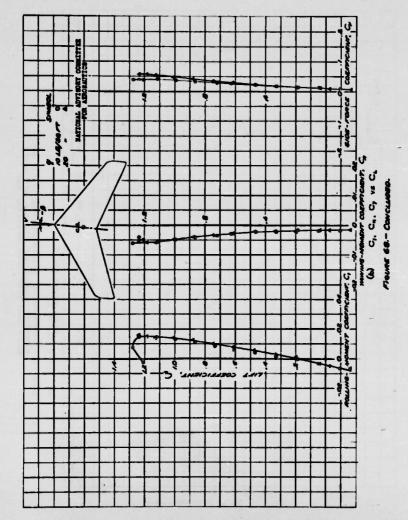
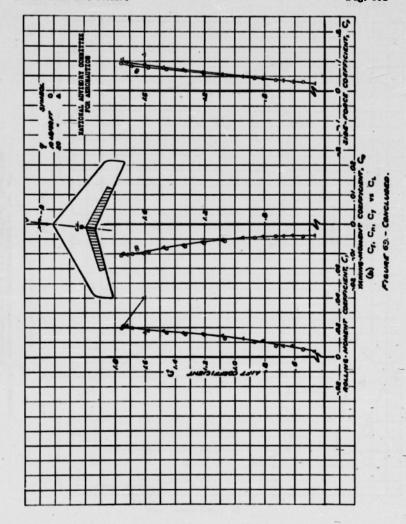


Fig. 68b



967 Fig. 69a NACA RM No. A6K15

Fig. 69b



98/

Fig. 70a

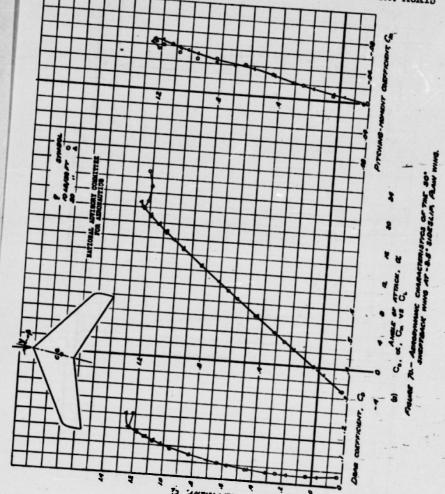
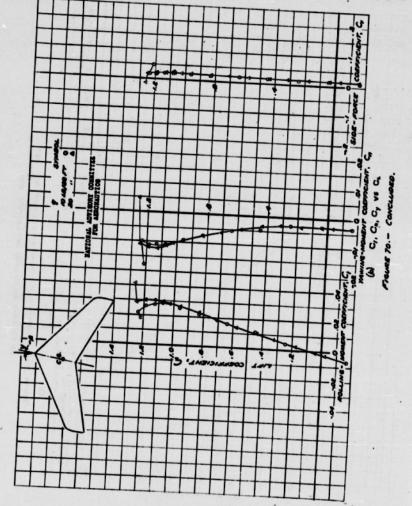
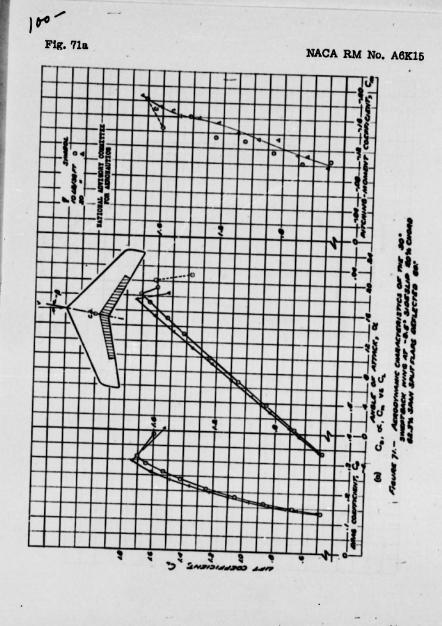


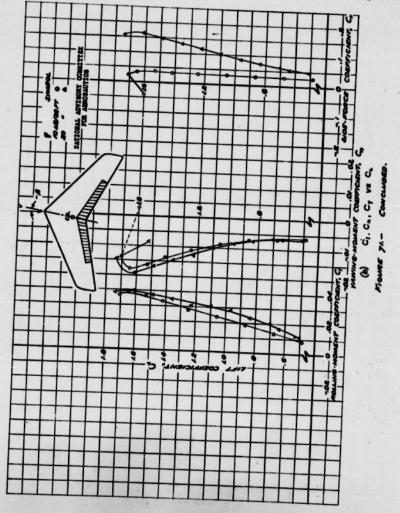
Fig. 70b





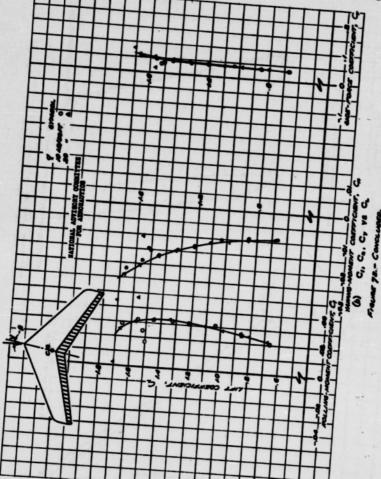
NACA RM No. A6K15





102 NACA RM No. A6K15 Fig. 72a

Fig. 72b

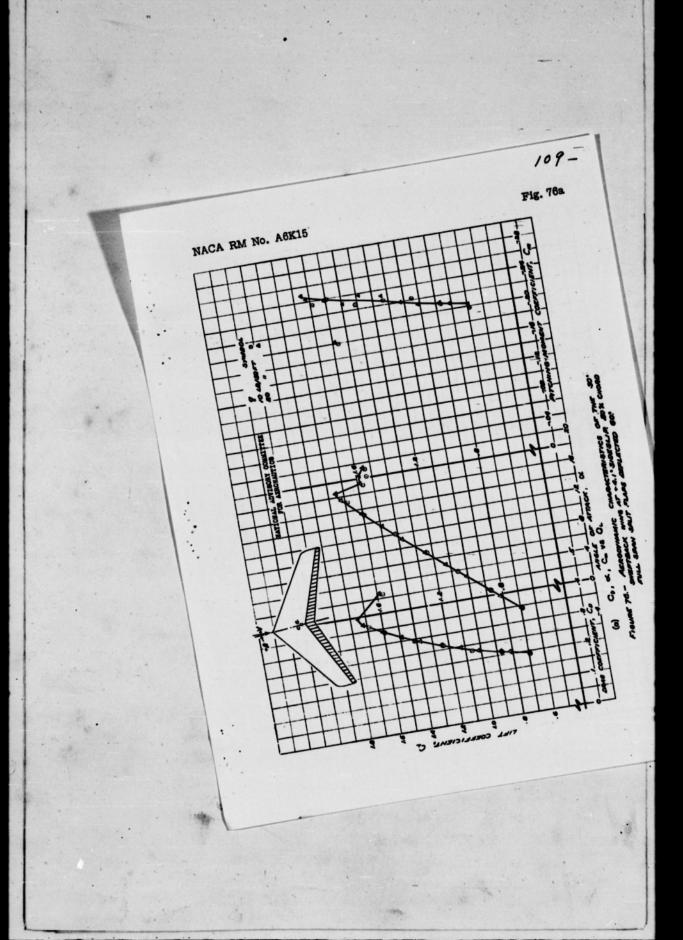


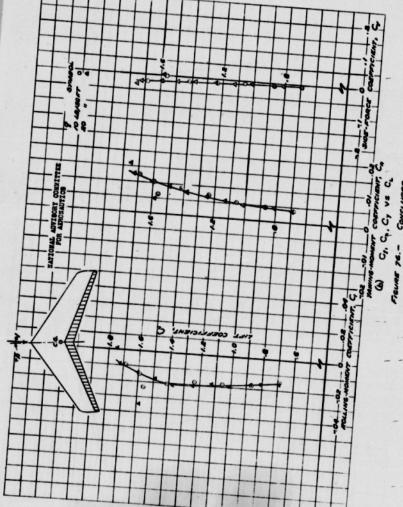
104 -Fig. 73 NACA RM No. A6K15 FIGURE 73. - Acrosy

105-Fig. 74a NACA RM No. A6K15

106 / Fig. 74b NACA RM No. A6K15

107-Fig. 75a NACA RM No. A6K15





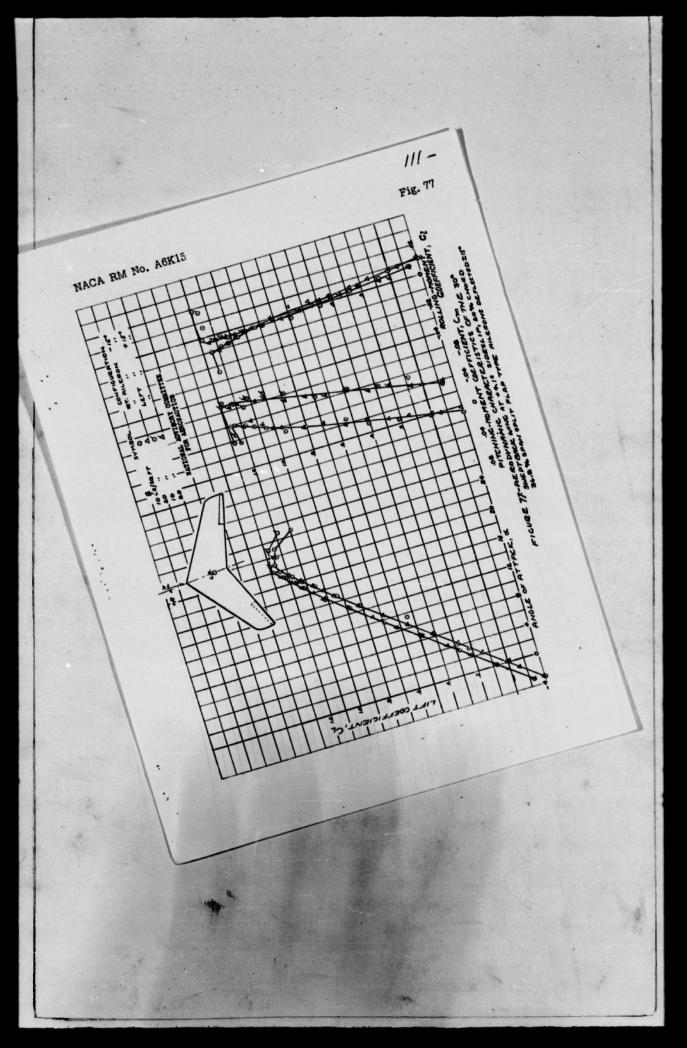
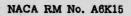
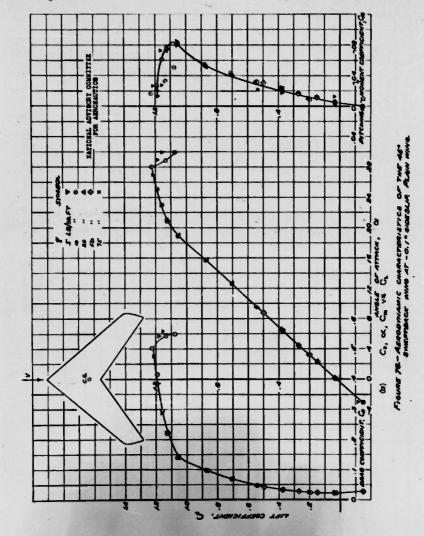


Fig. 78a





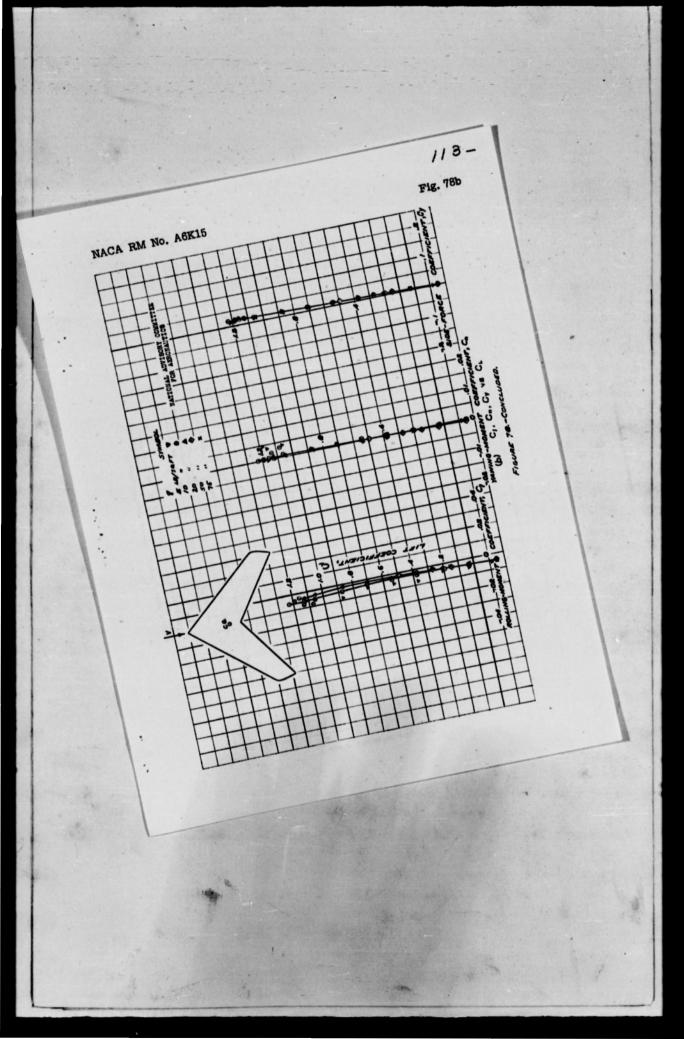
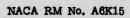
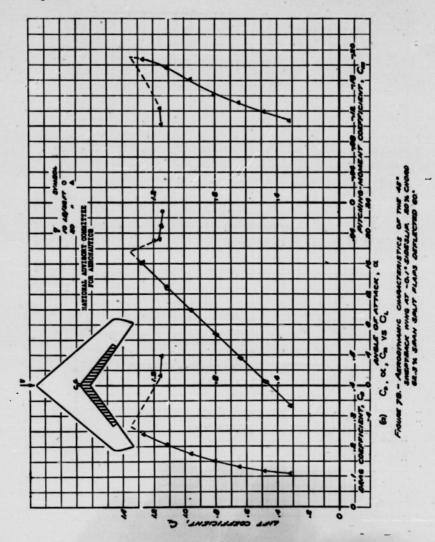
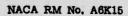


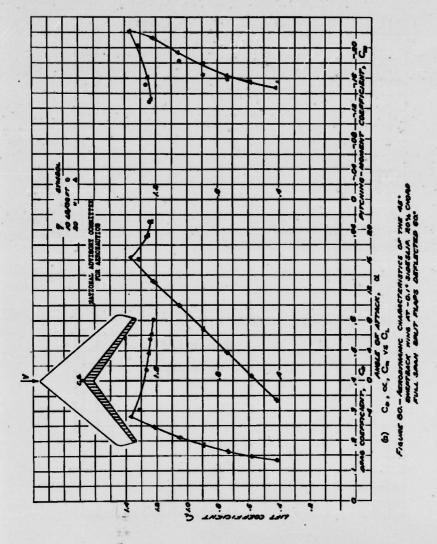
Fig. 79a





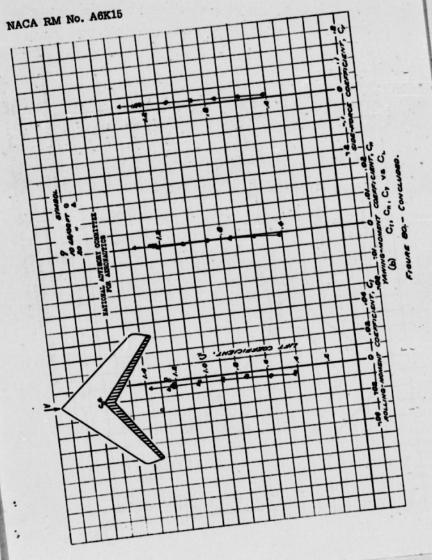
115-Fig. 79b NACA RM No. A6K15

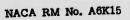


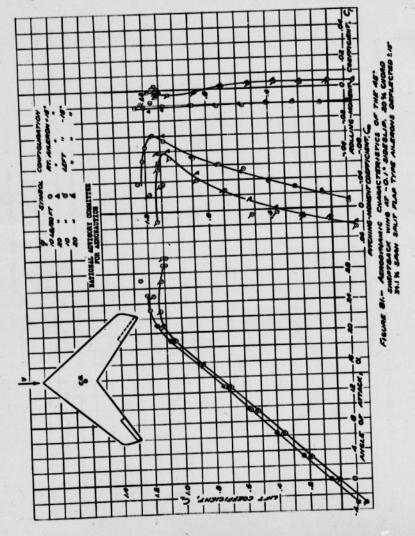


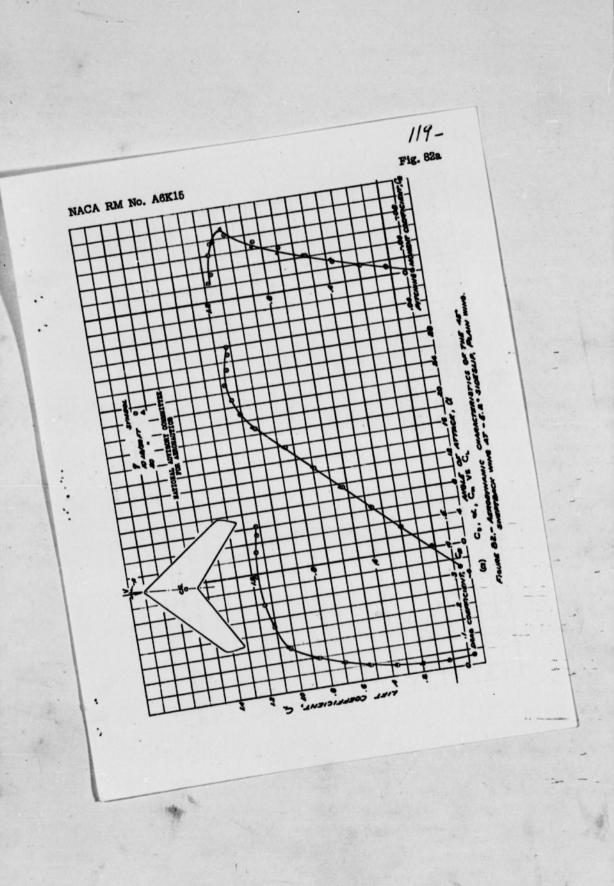
117-

Fig. 80b





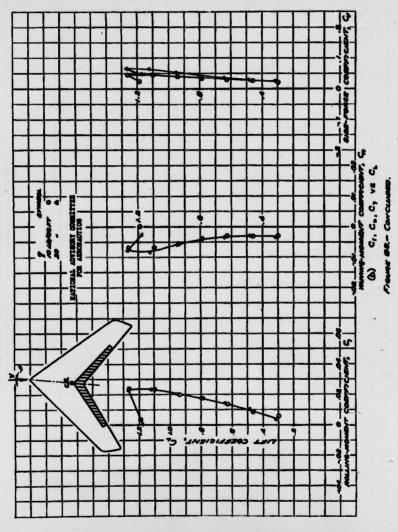


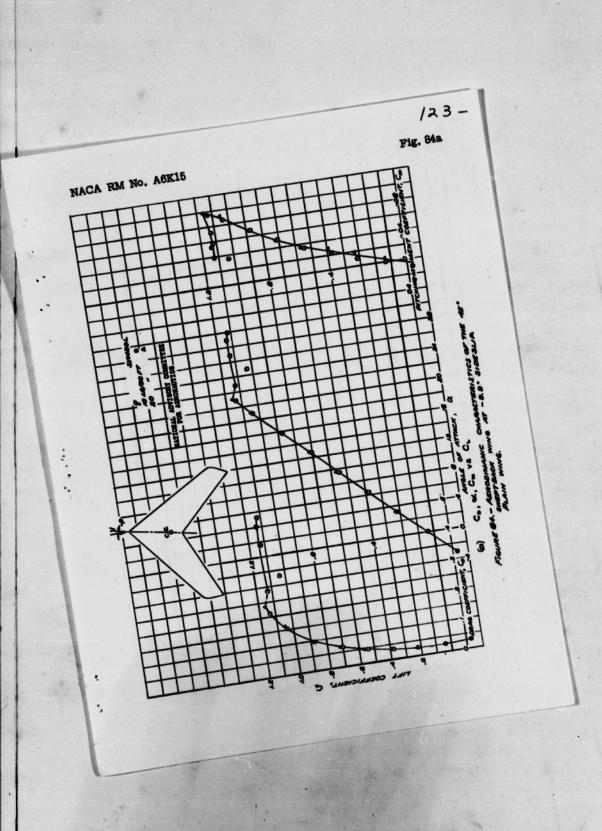


121-Fig. 83a NACA RM No. A6K15

122

Fig. 83b

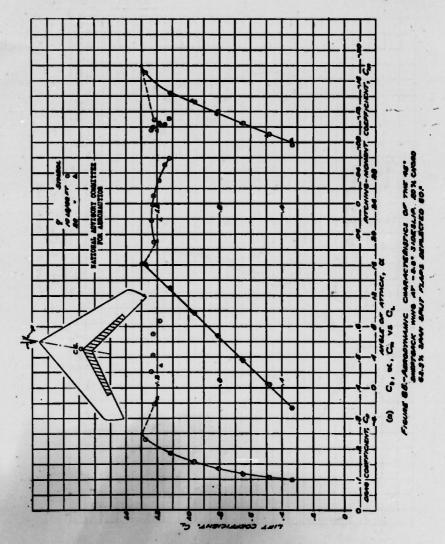




124-NACA RM No. A6K15 Fig. 84b

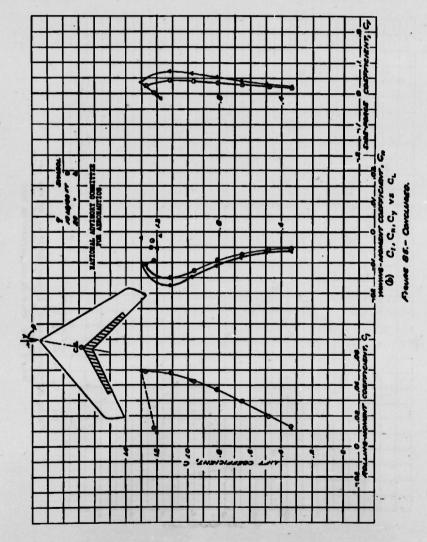
NACA RM No. A6K15



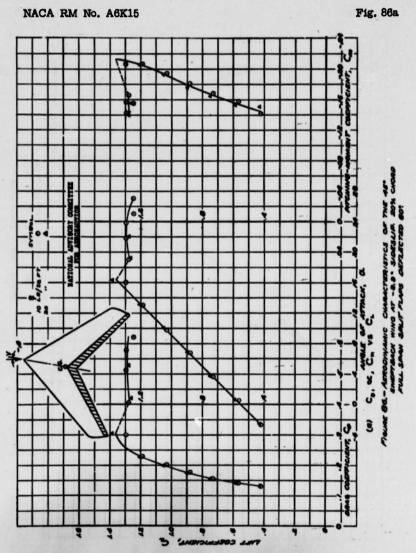


126

Fig. 85b

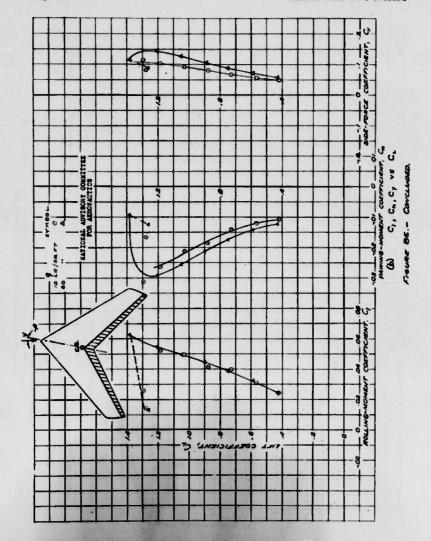


NACA RM No. A6K15

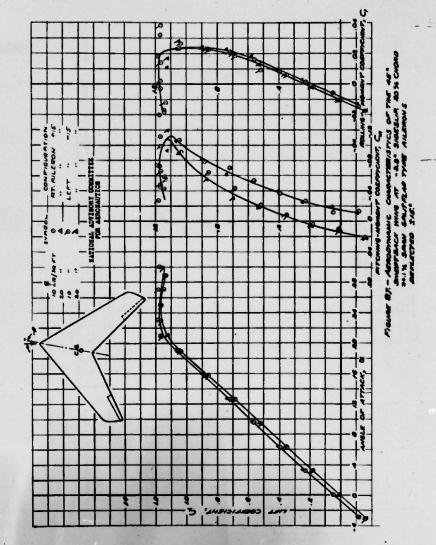


1281

Fig. 86b







Pig. 88a

NACA RM No. A6K15

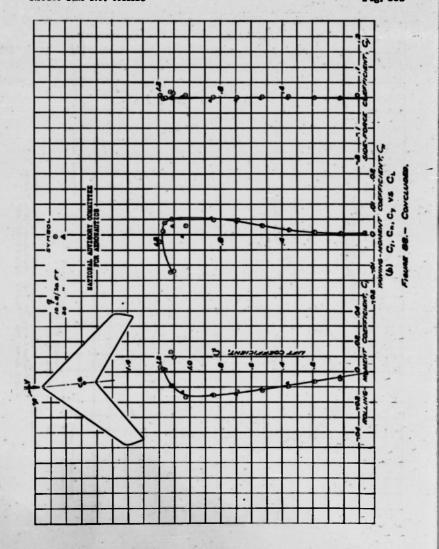
Pig. 88a

(9)

NACA RM No. A6K16

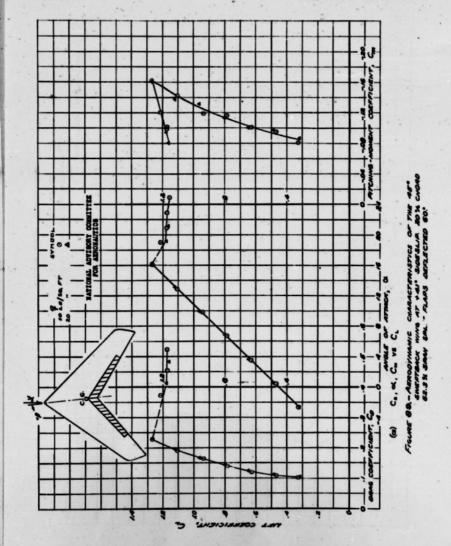
NACA RM No. A6K15





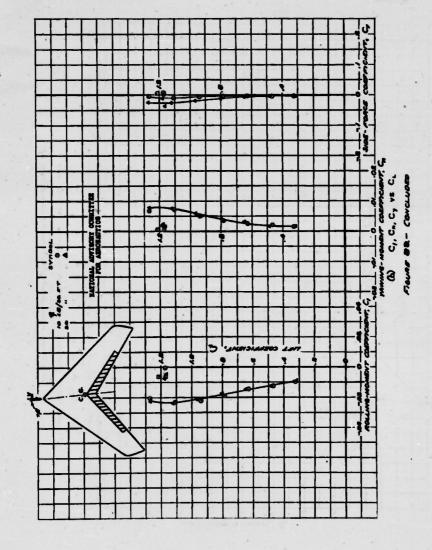
132

Fig. 89a



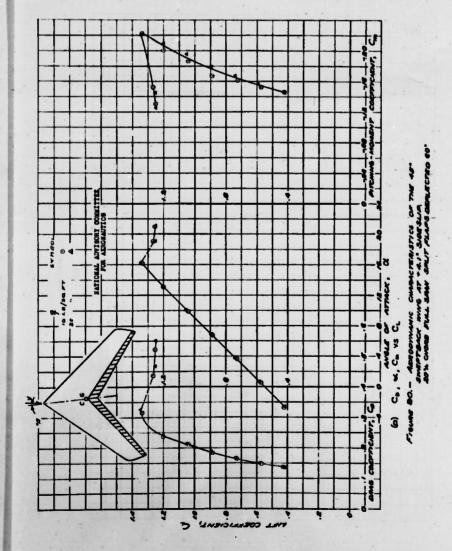
NACA RM No. A6K15

Fig. 89b



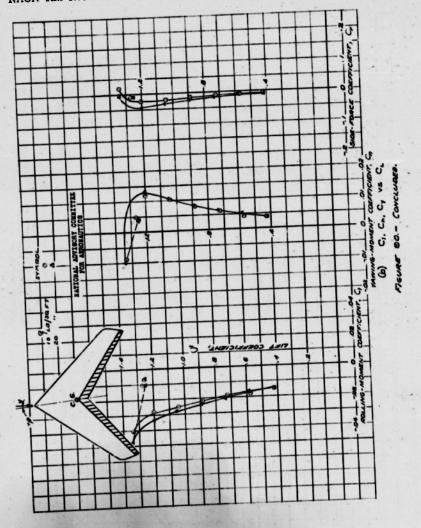
134-

Fig. 90a



135-

Fig. 90b



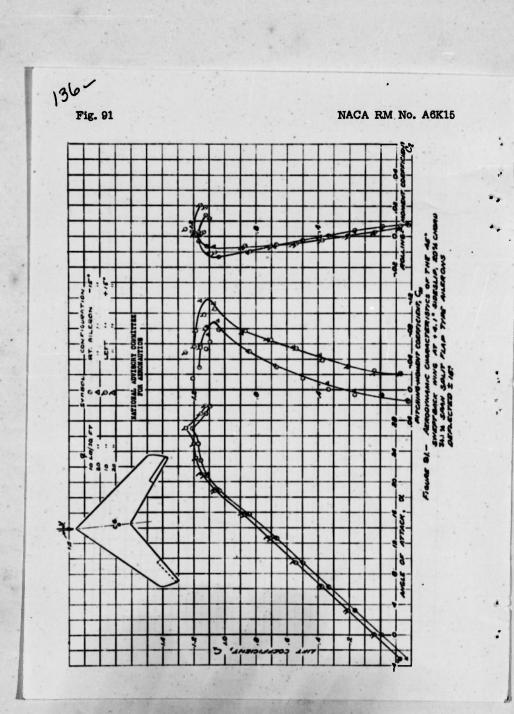


Fig. 92

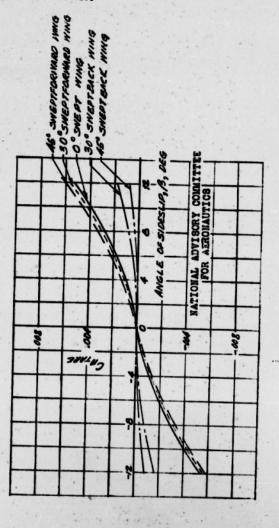


FIGURE 92.-YANING-MINENT COEFFICIENT TARE OF THE FIVE FIVE SWEET WINGS. (ADMILED TO BATIC CANA TO COTAIN SUMMARY UNTAL)

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